

**2002 Agilent Technologies
Europhysics Prize Lecture
on**

Quantum Dynamics of Nanomagnets

Bernard Barbara, *L. Néel Lab, Grenoble, France*

Jonathan R. Friedman, *Amherst College, Amherst, MA, USA*

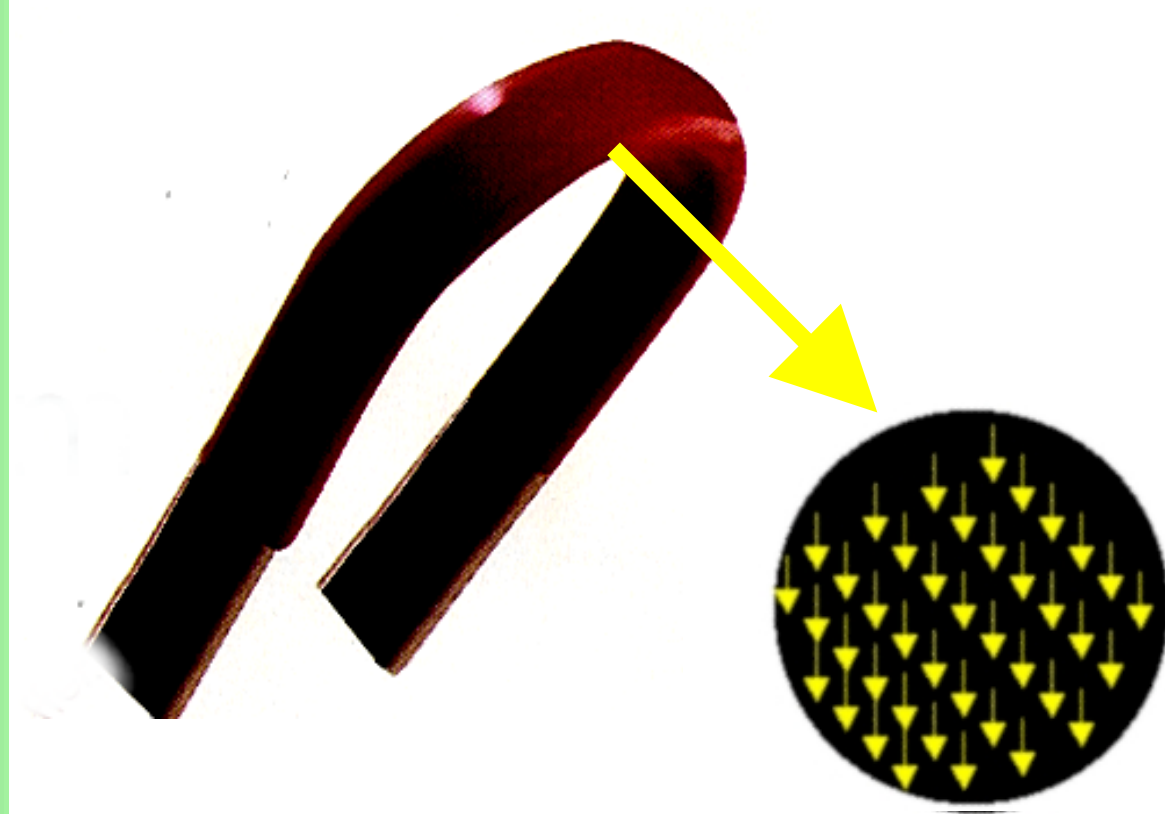
Dante Gatteschi, *University of Florence, Italy*

Roberta Sessoli, *University of Florence, Italy*

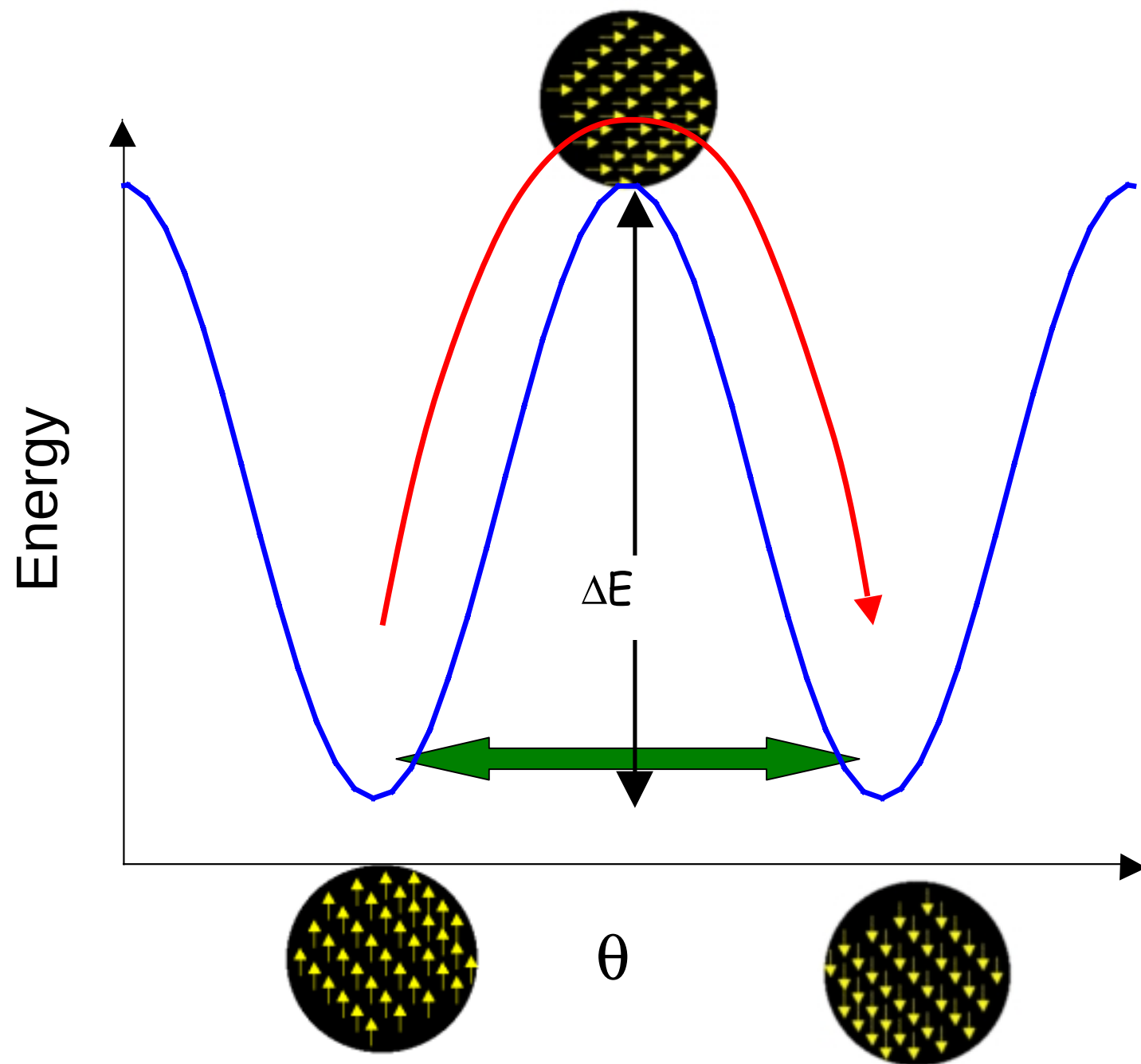
Wolfgang Wernsdorfer, *L. Néel Lab, Grenoble, France*

Budapest 26/08/2002

the miniaturization process



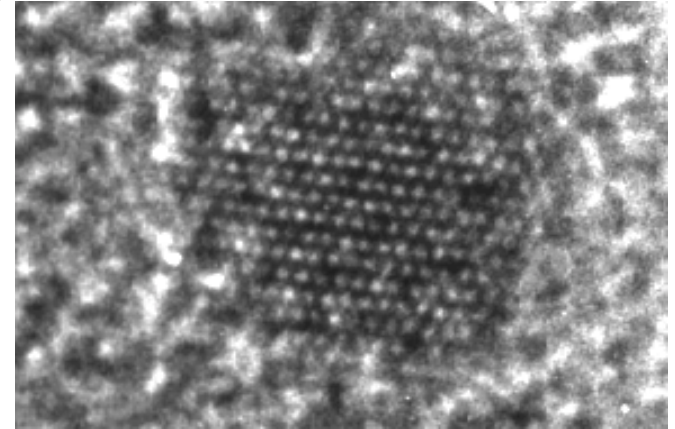
Single Domain Particles
coherent rotation of all the spins



Quantum effects in the dynamics of the magnetization

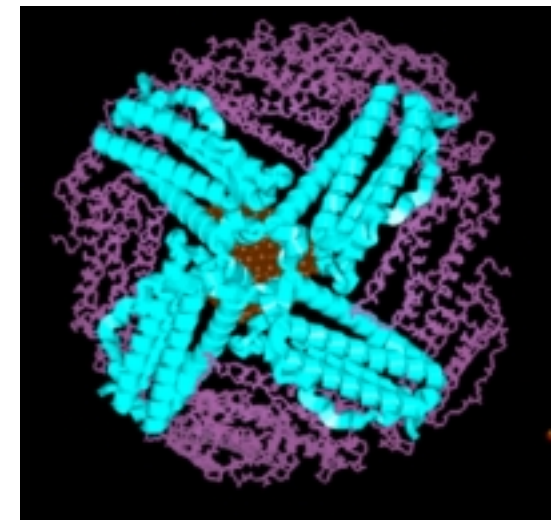
- First evidences of Quantum Tunneling in nanosized magnetic particles

(difficulties due to size distribution)



0 3 nm

- Quantum Coherence in ferrihydrite confined in the ferritin mammalian protein
- (inconclusive due to distribution of iron load)





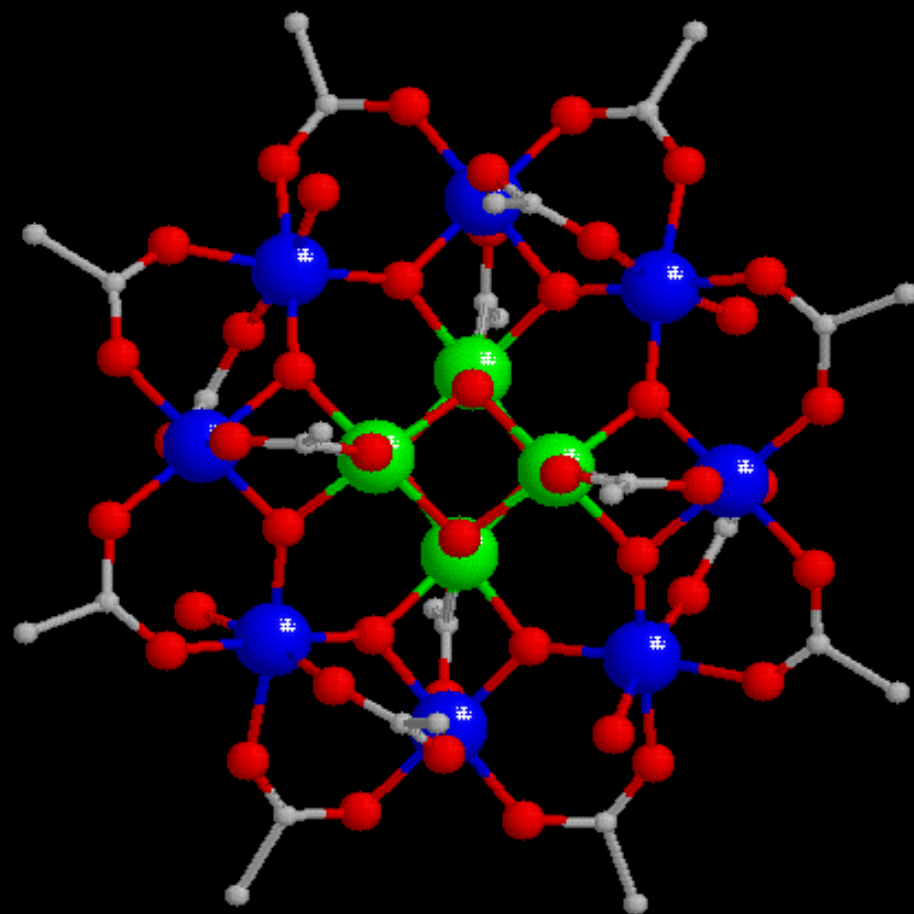
= metal ions

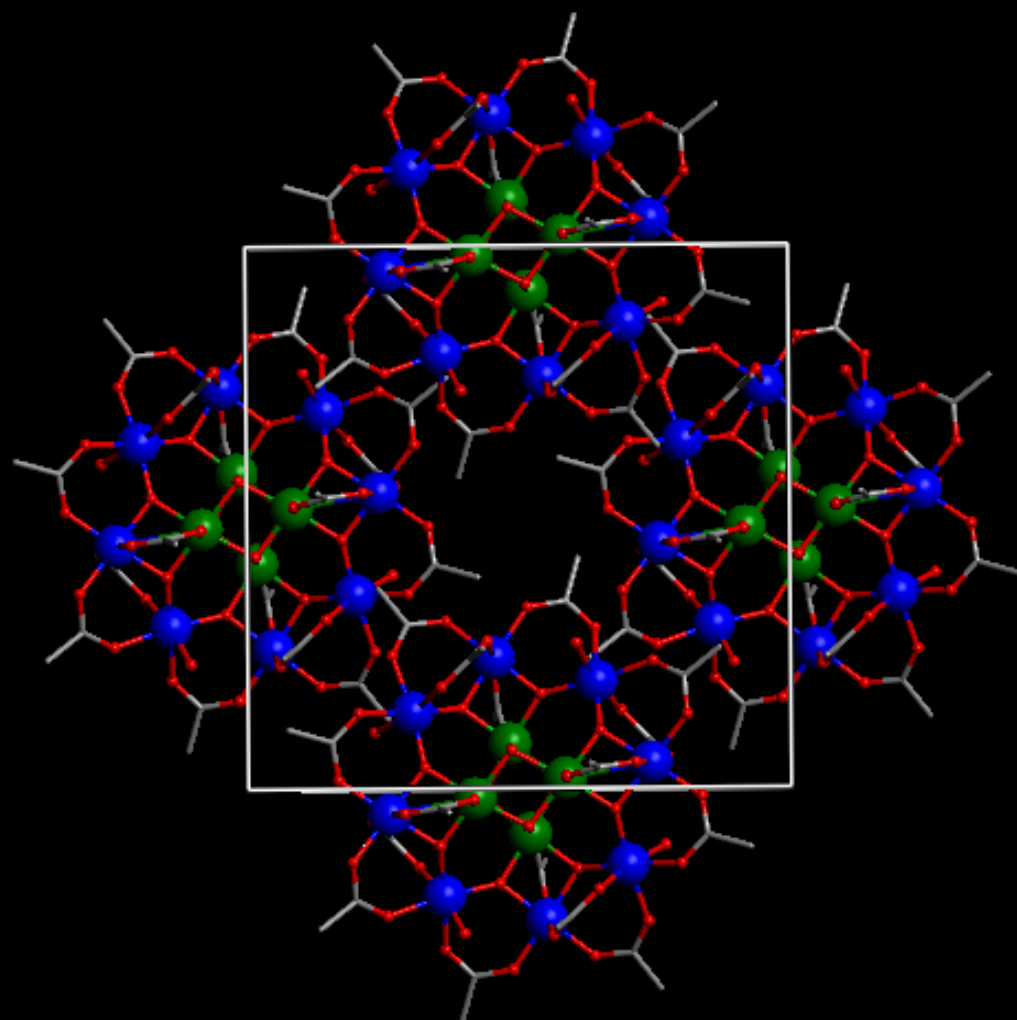


= oxygen

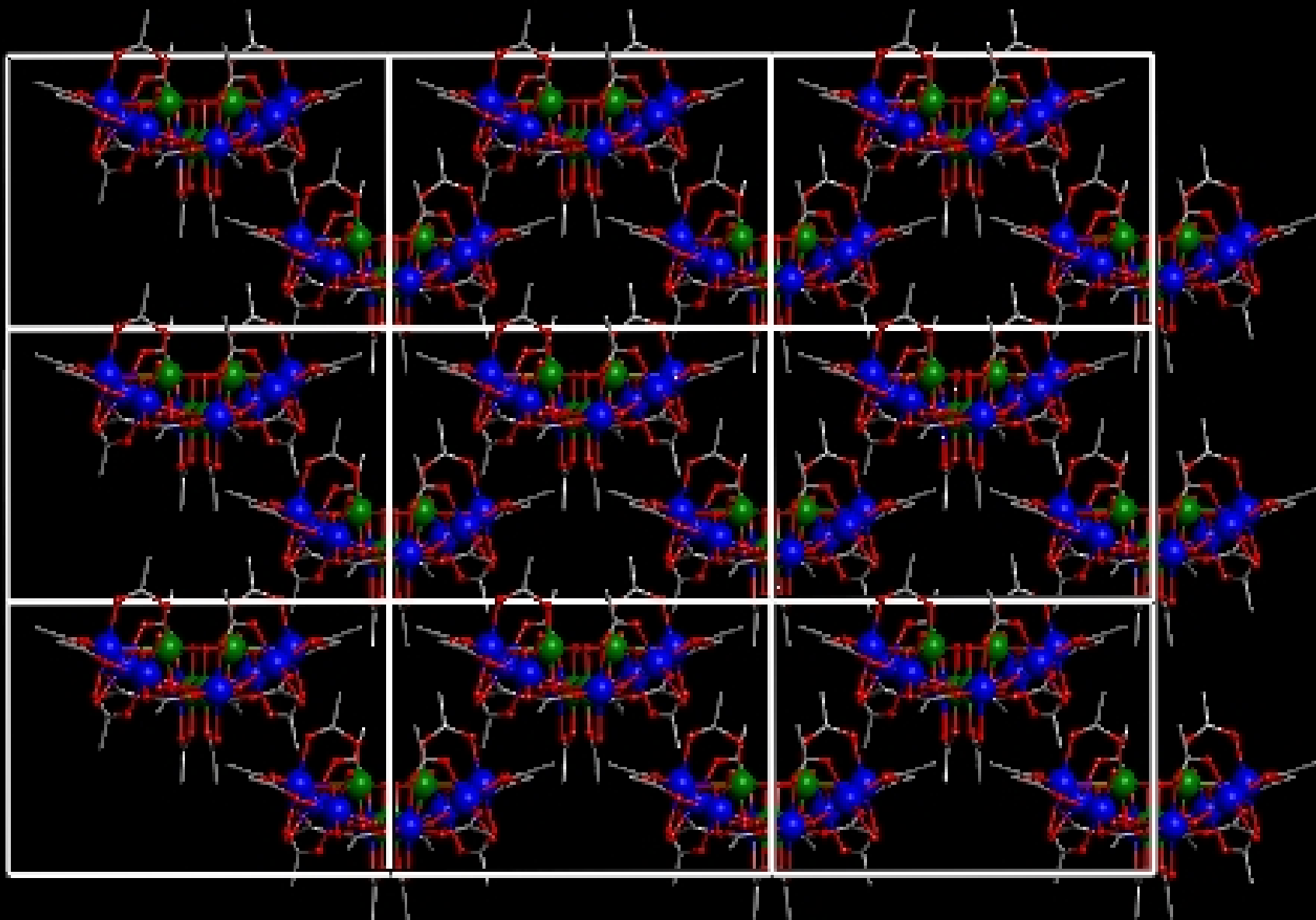


= carbon





The molecules are regularly arranged in the crystal

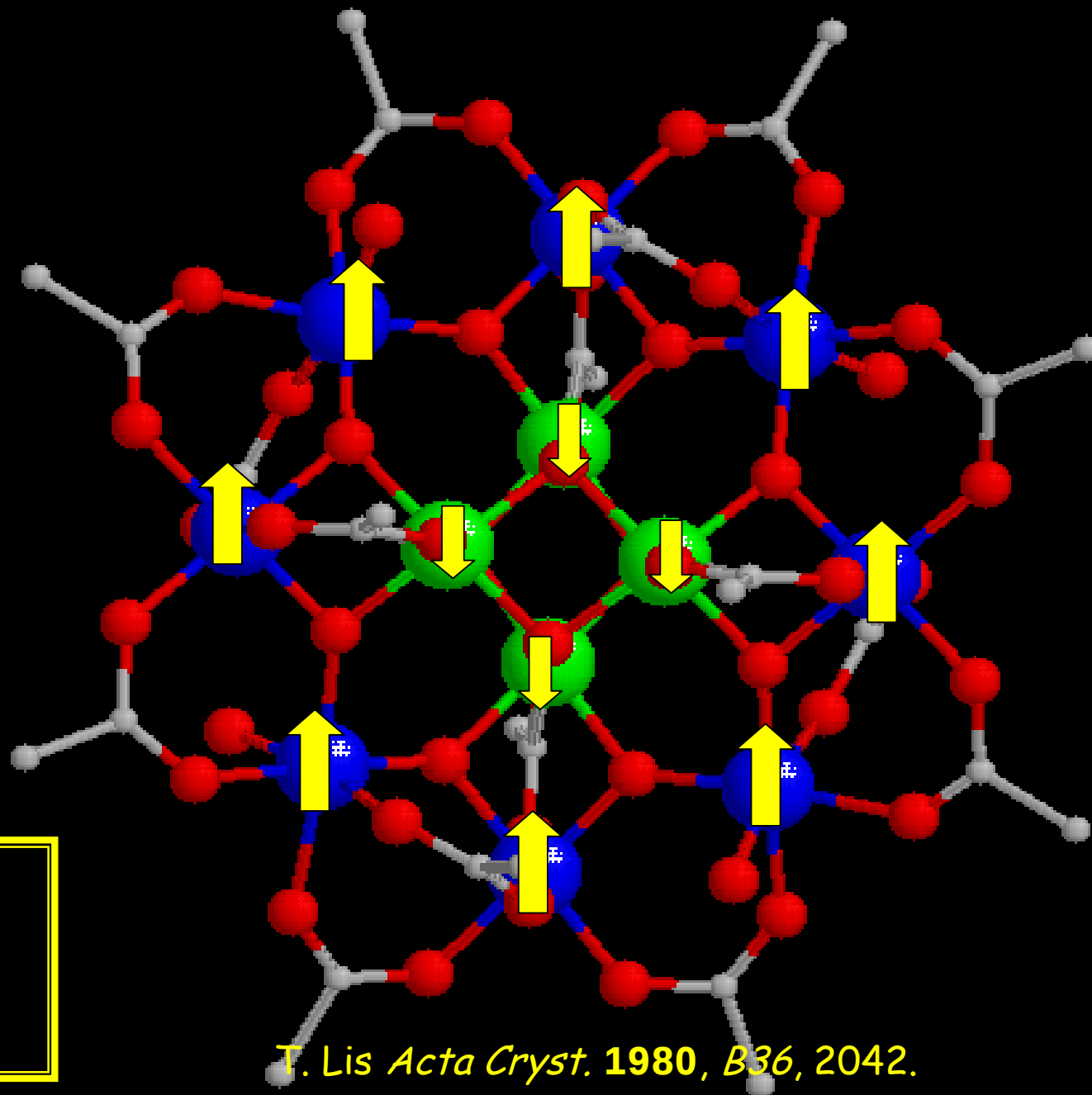


Mn12acetate

● Mn(III)
 $S=2$

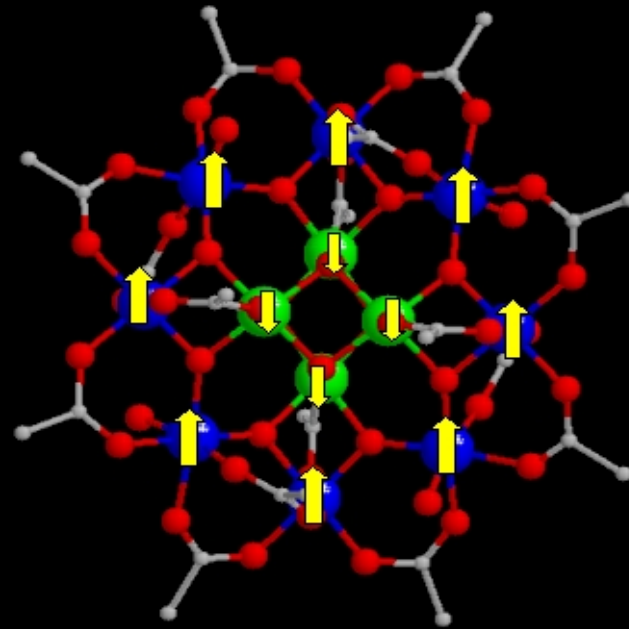
● Mn(IV)
 $S=3/2$

Total Spin
 $=10$

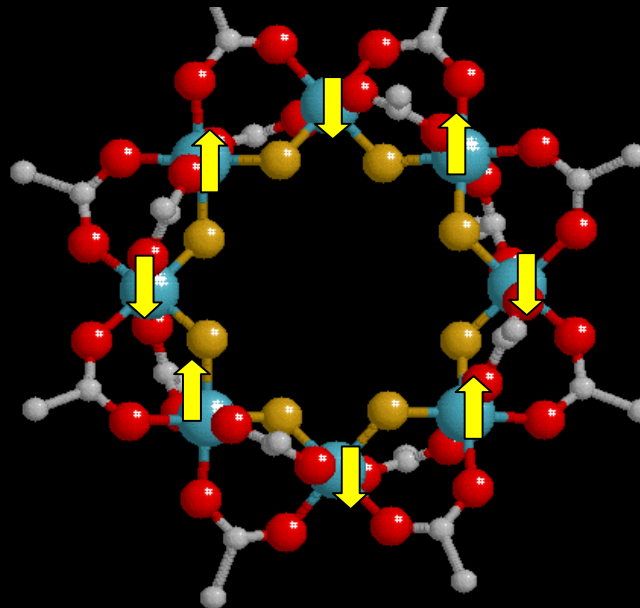


T. Lis *Acta Cryst.* **1980**, B36, 2042.

high spin molecules



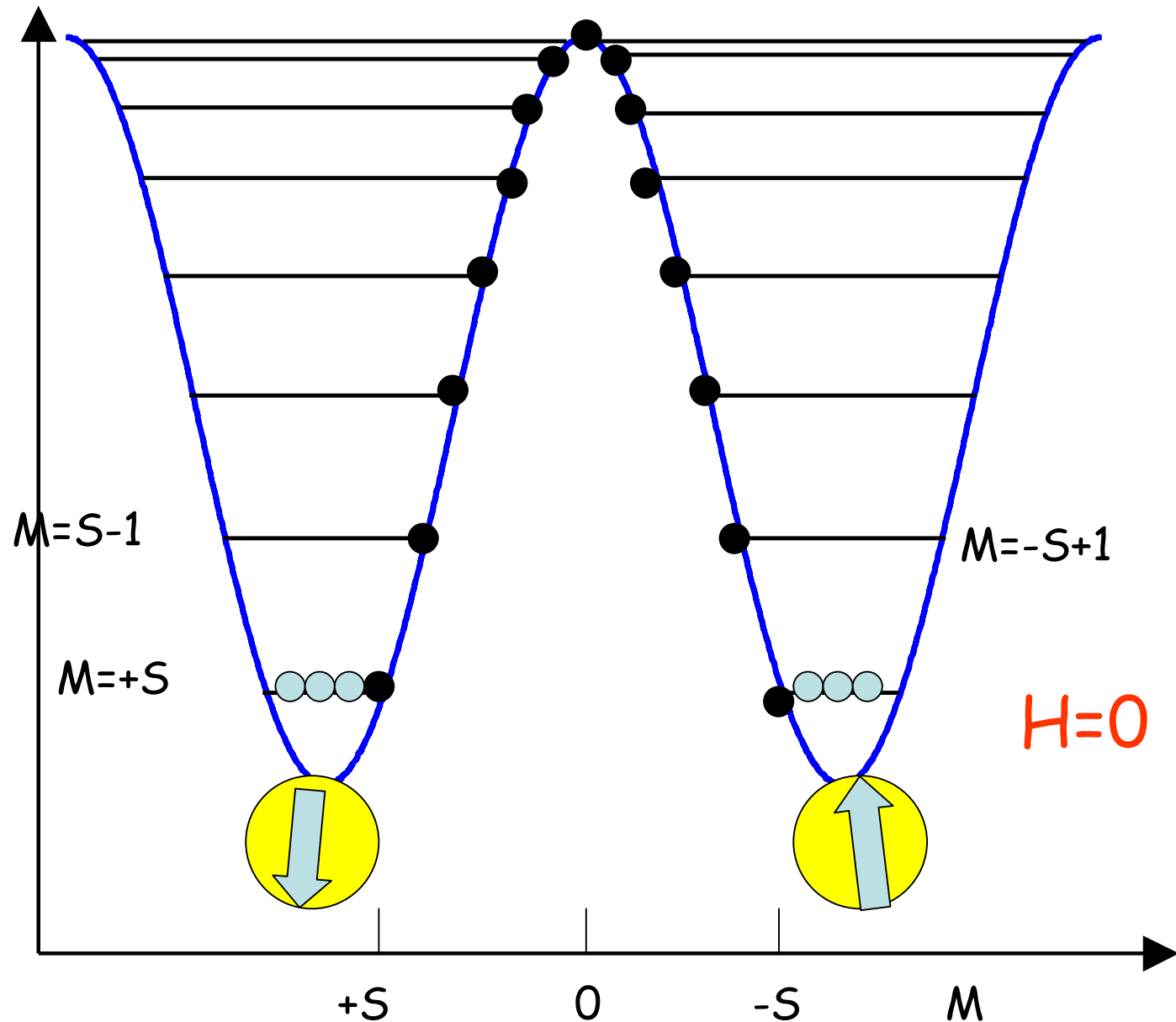
and low spin molecules



Uniaxial magnetic anisotropy

$$H = -DS_z^2$$

If S is large

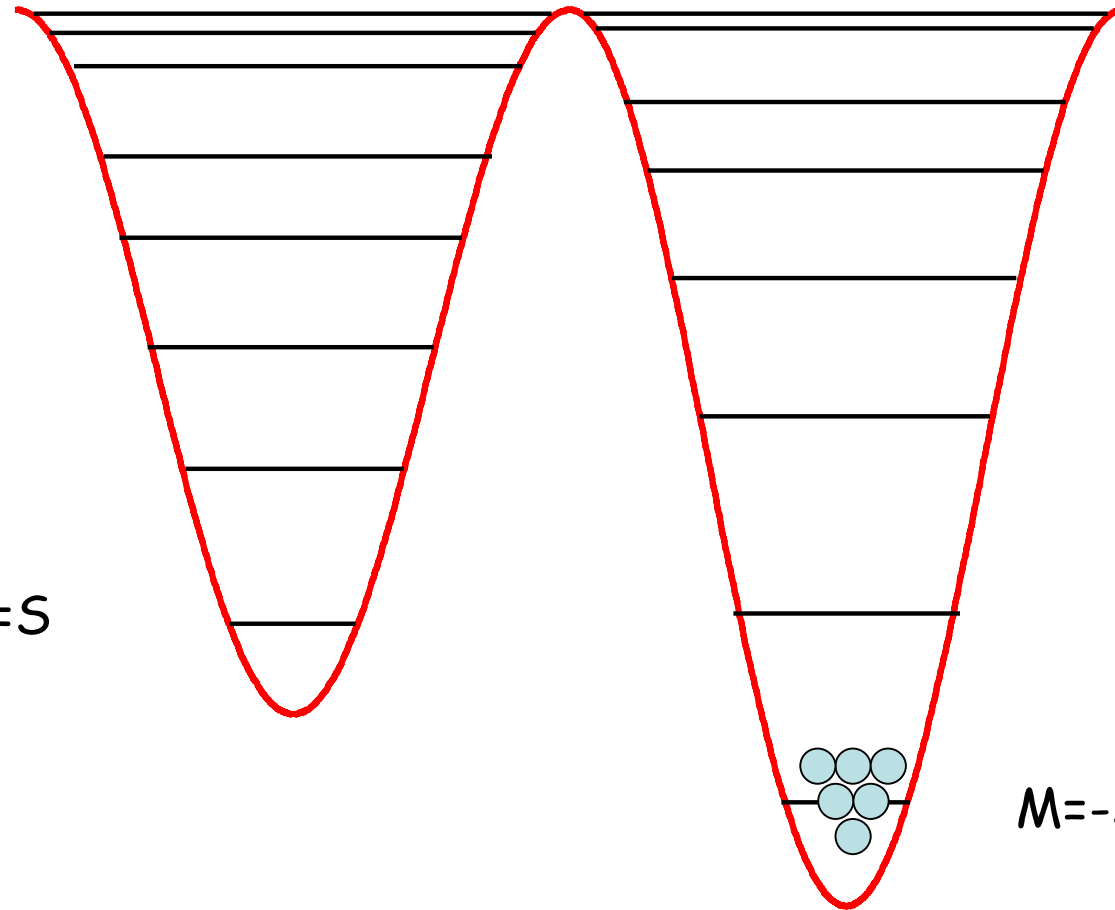


$$H = -DS_z^2 + g\mu_B H_z S_z$$

$H \neq 0$

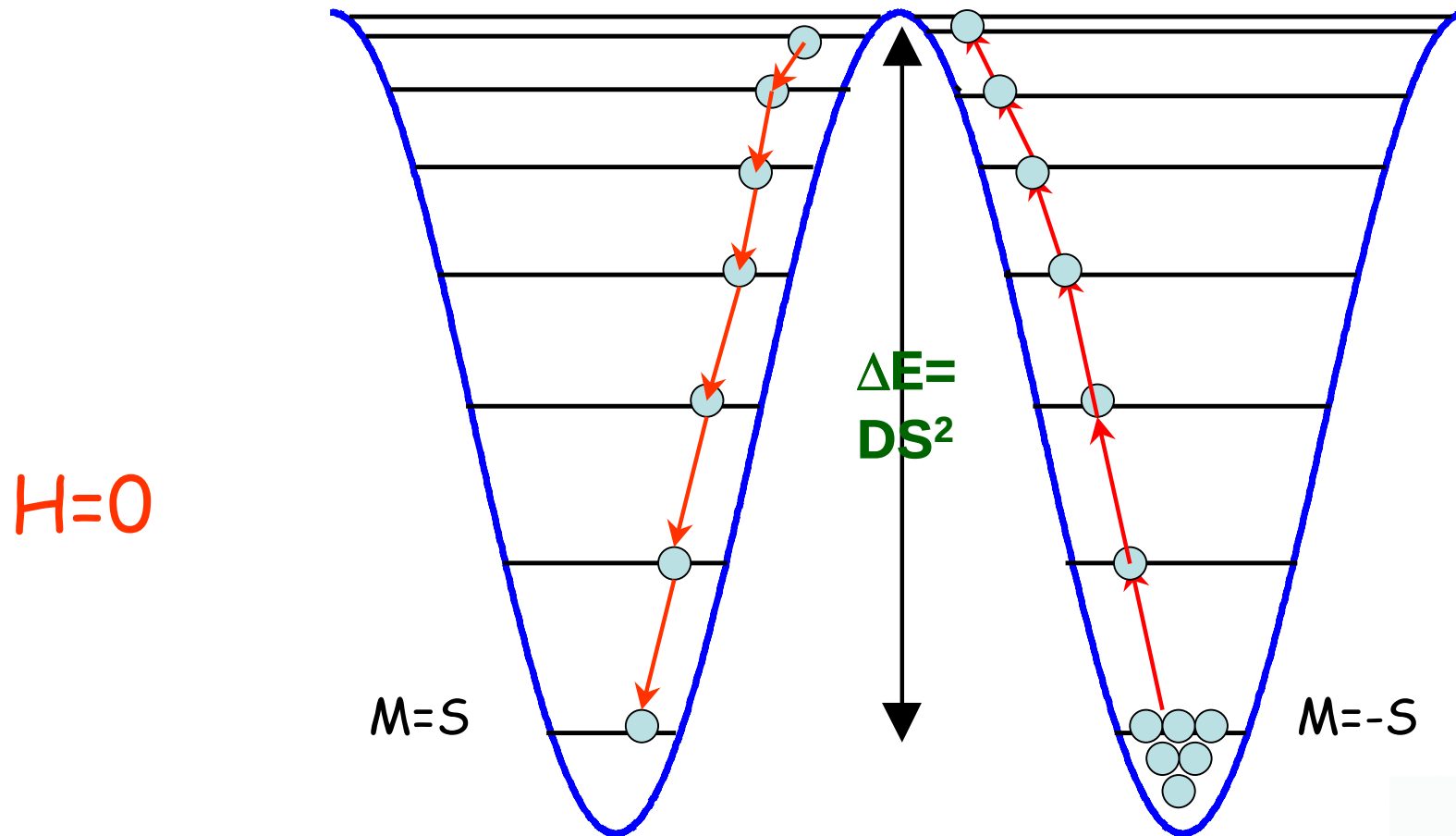
$M = S$

$M = -S$



return to the equilibrium

thermal activated mechanism

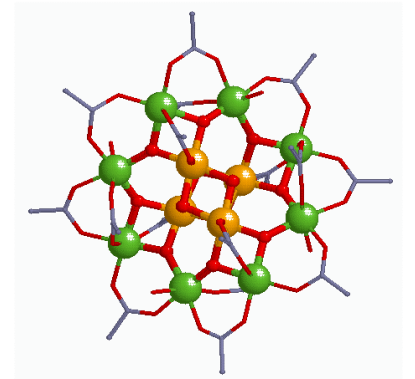


$$\tau = \tau_0 \exp(\Delta E / k_B T)$$

time=0

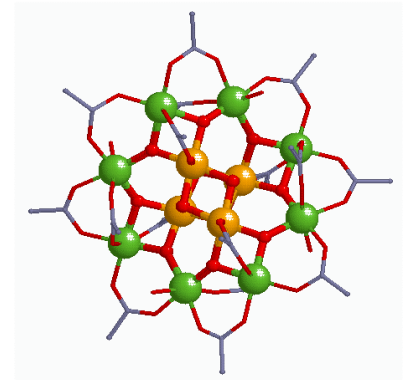
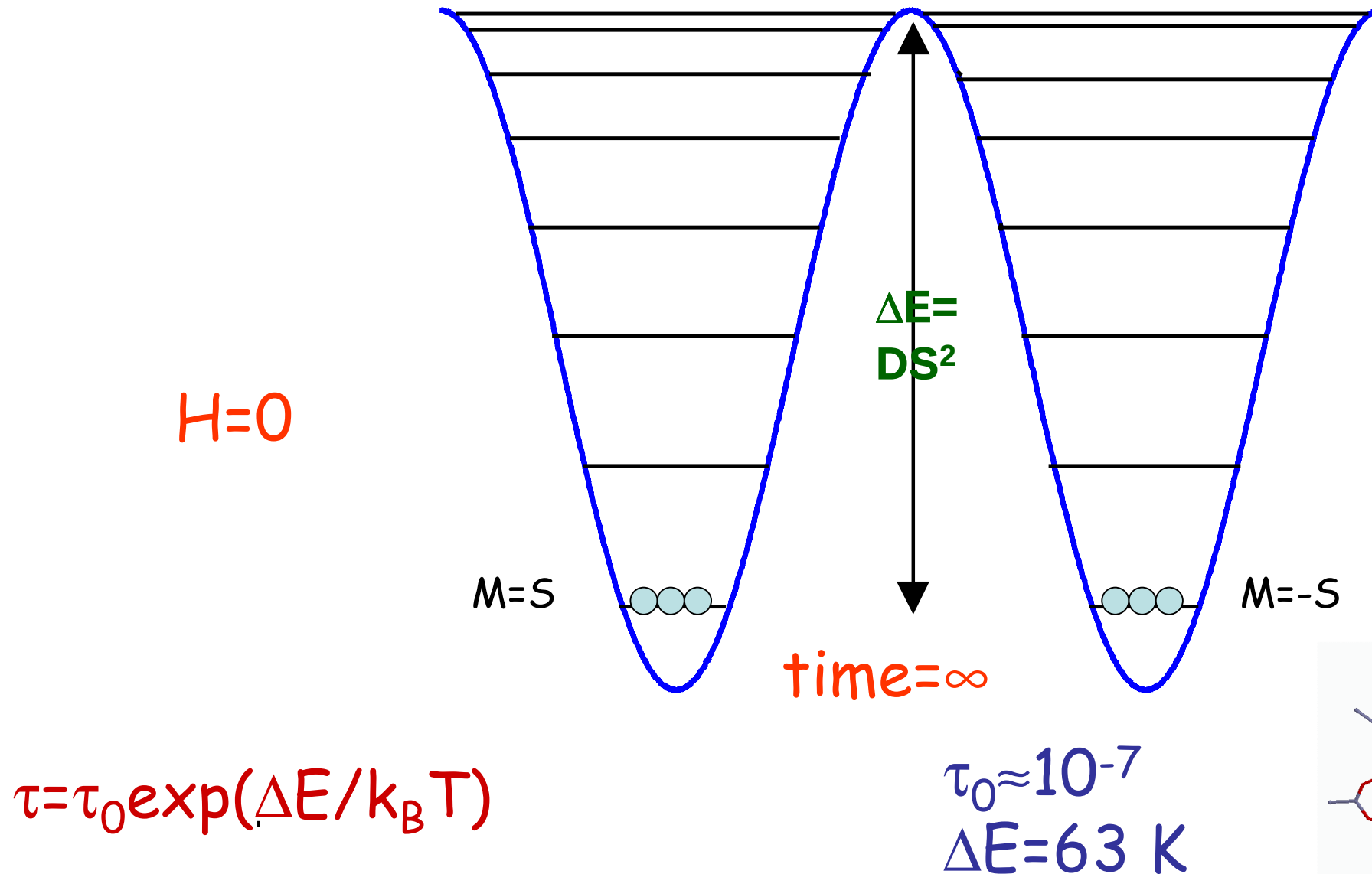
$$\tau_0 \approx 10^{-7}$$

$$\Delta E = 63 \text{ K}$$

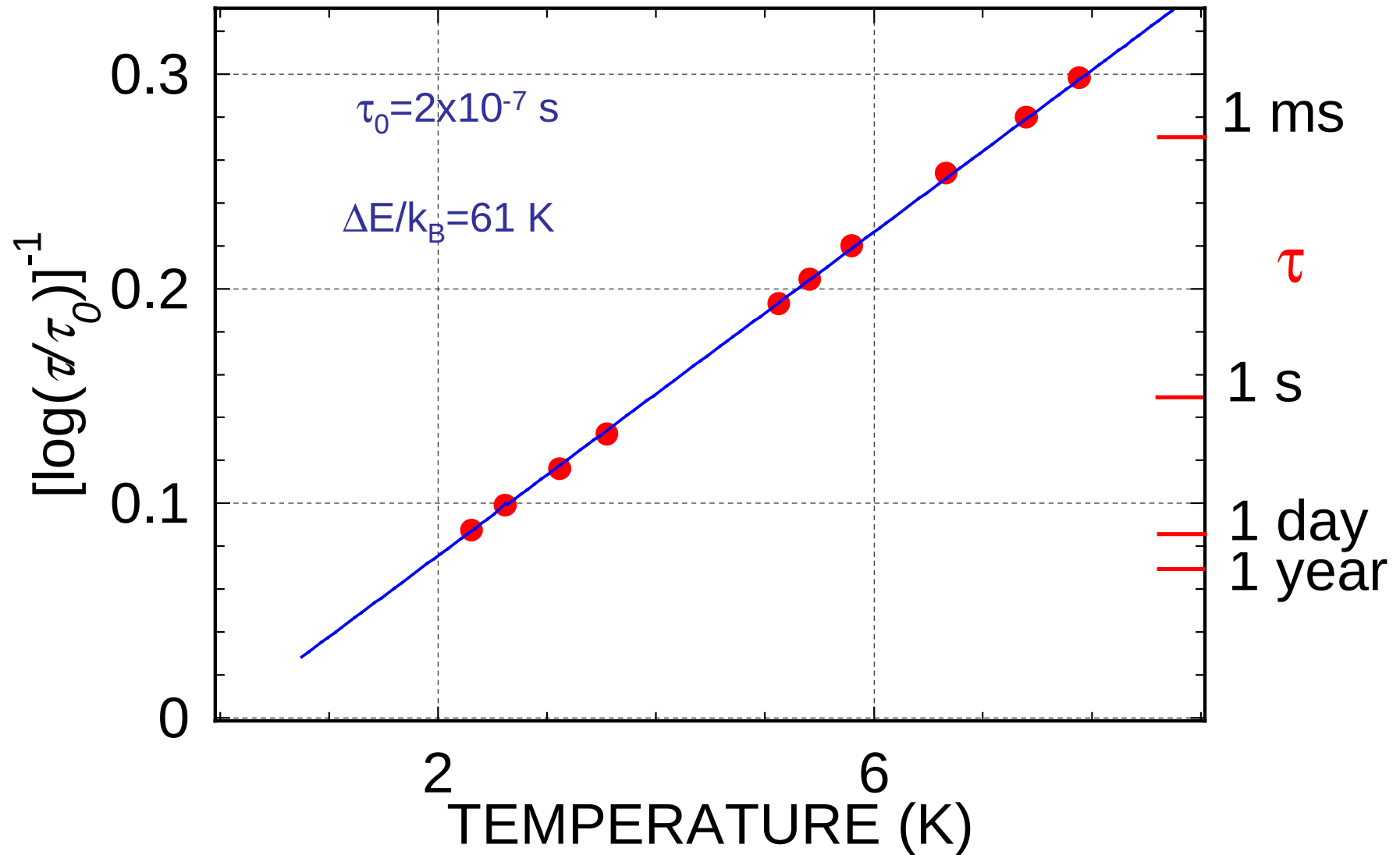


return to the equilibrium

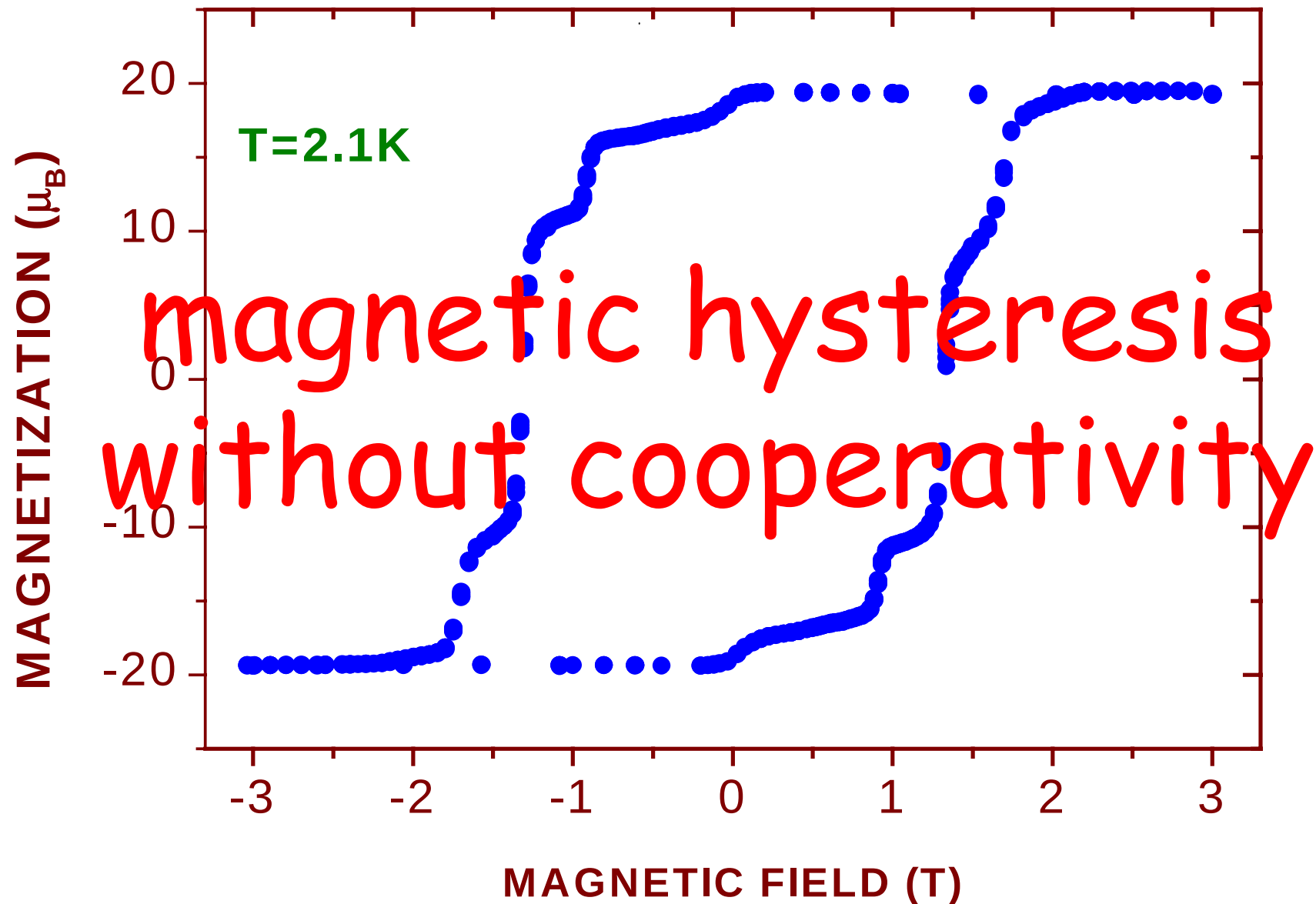
thermal activated mechanism



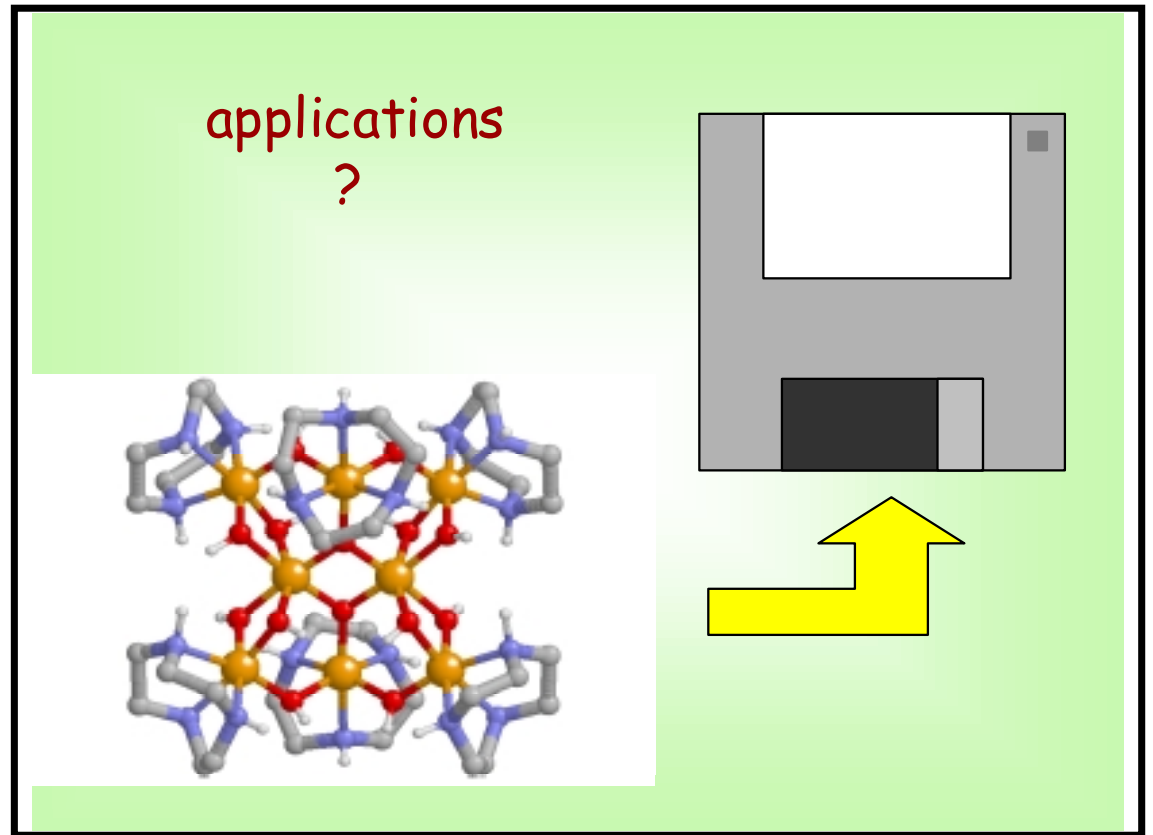
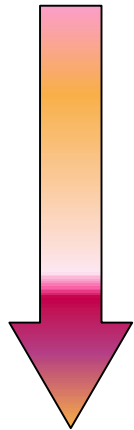
Temperature dependence of the relaxation time of Mn12acetate



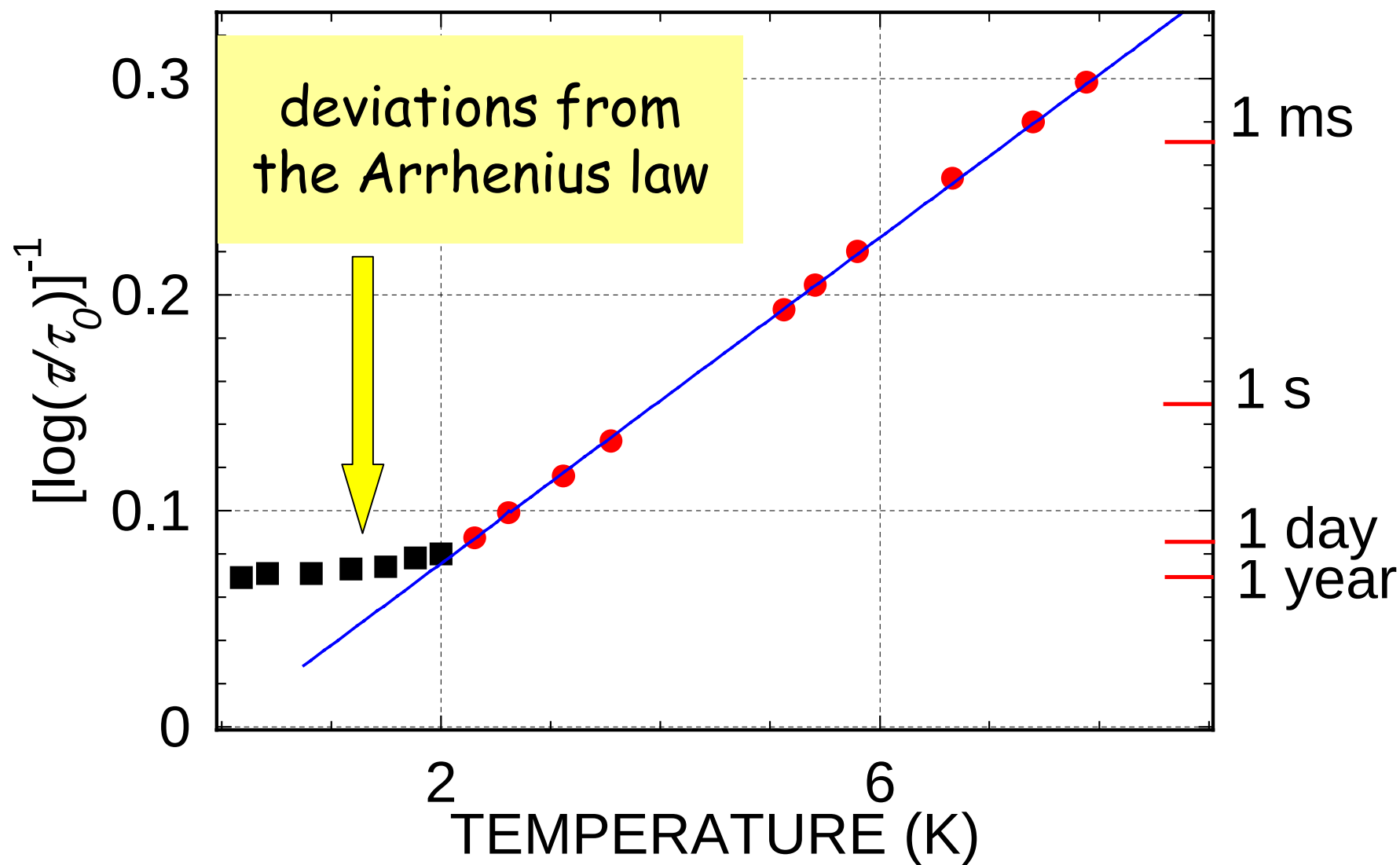
Mn12acetate: Hysteresis loop



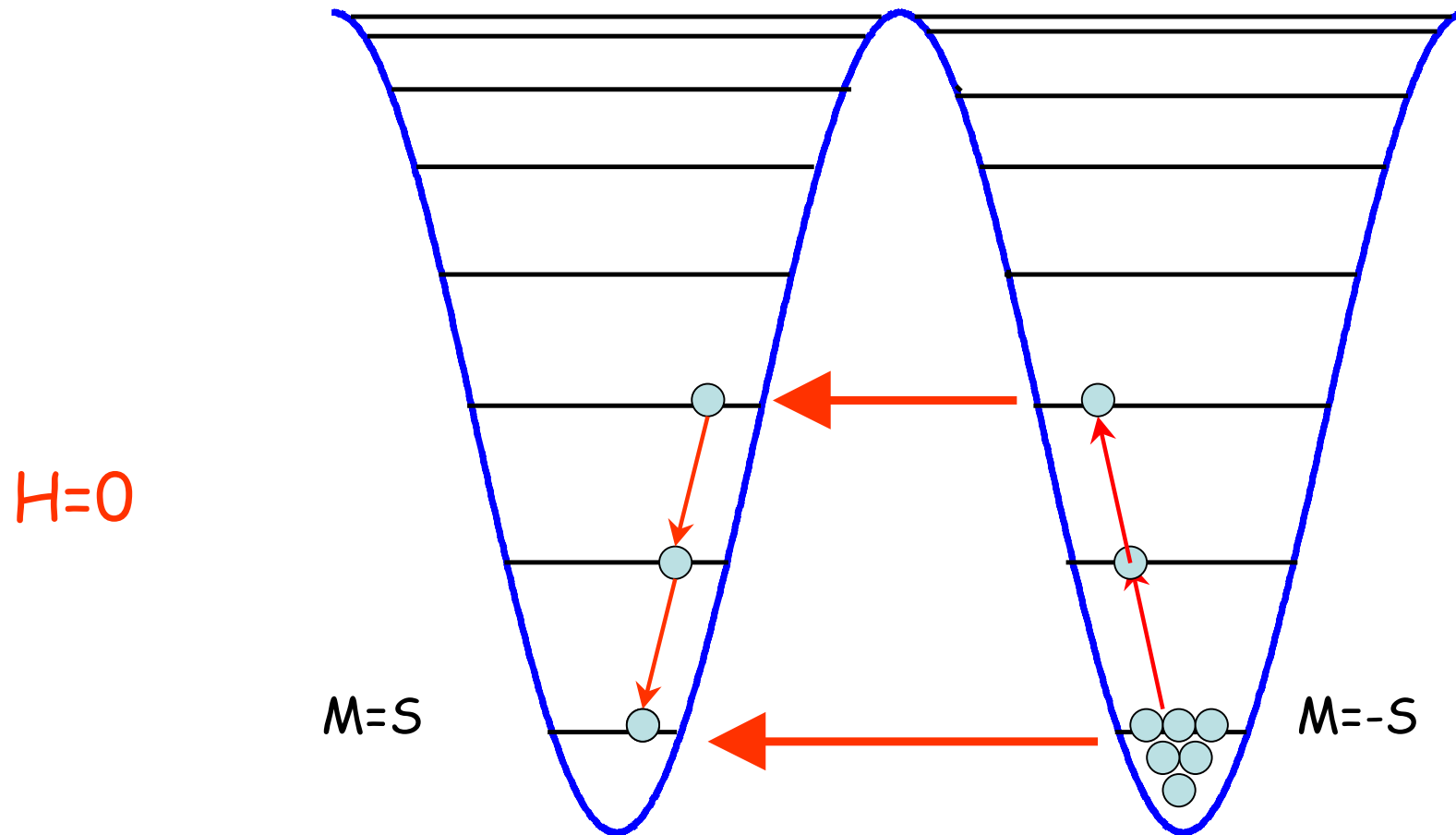
High Spin Clusters



Single Molecule Magnets

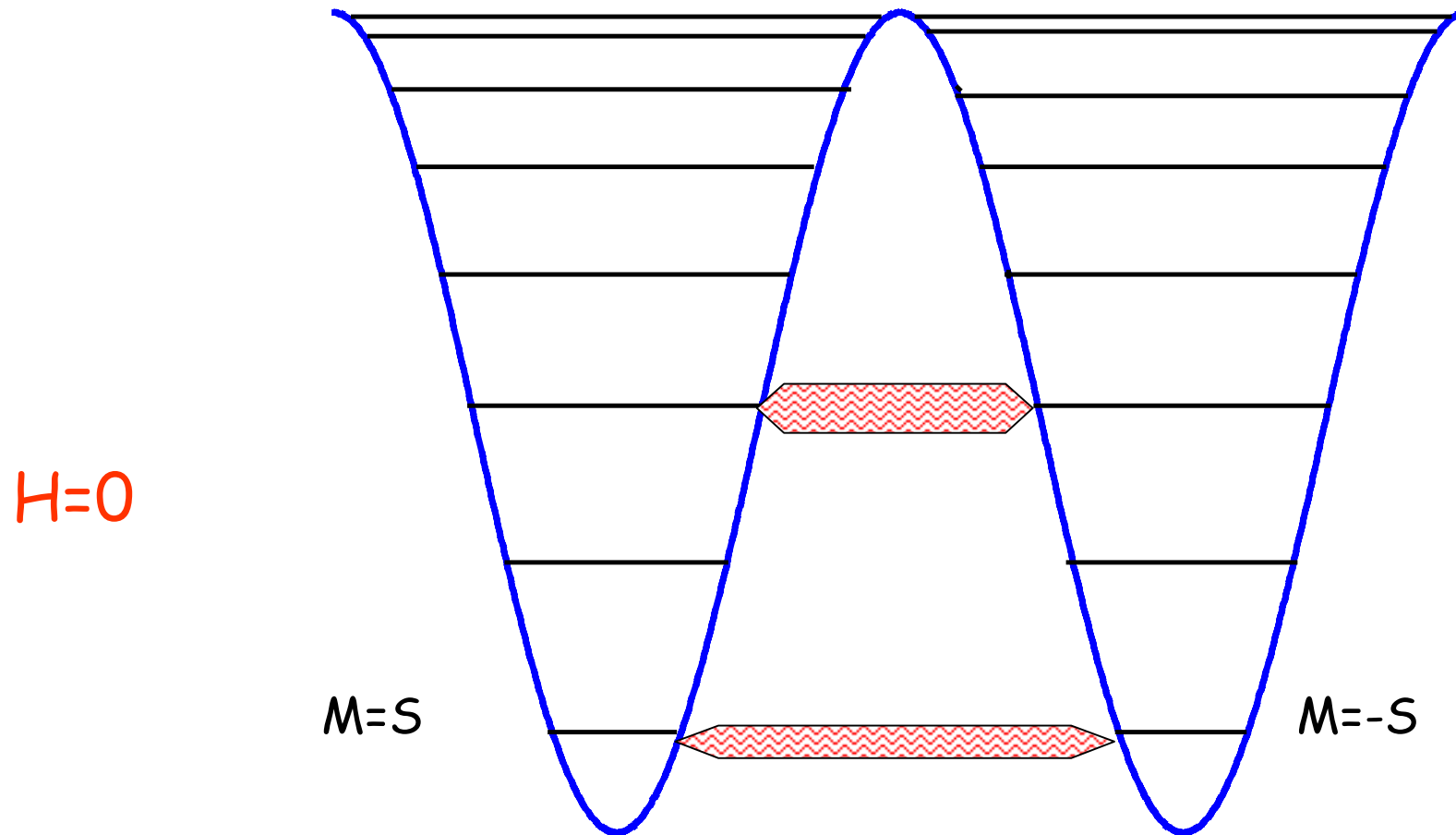


return to the equilibrium tunnel mechanism



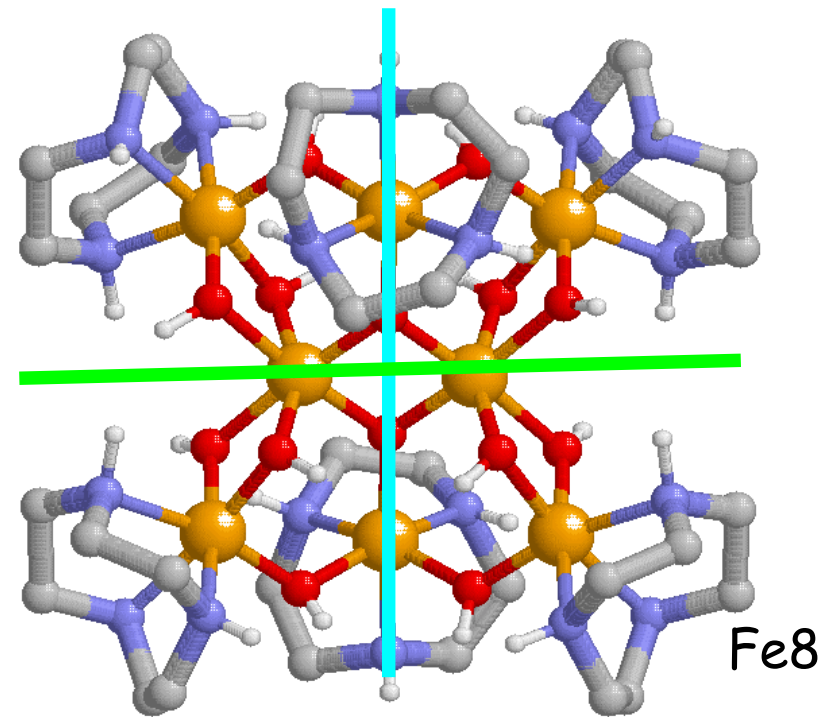
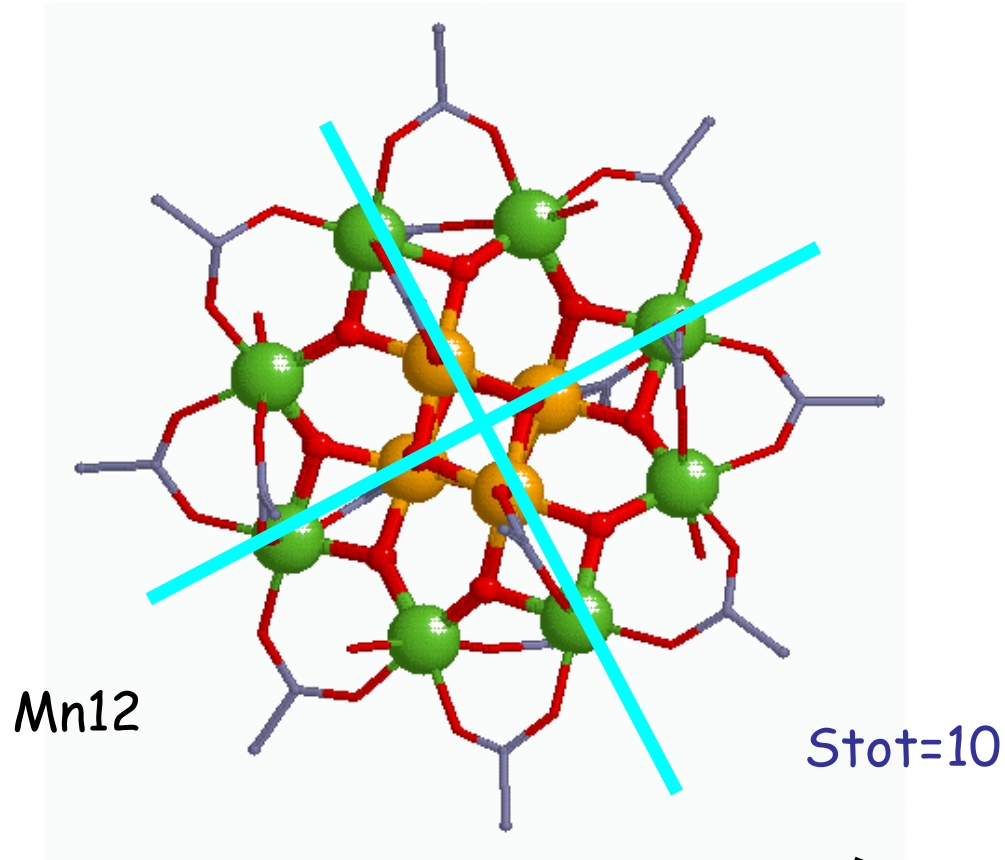
terms in S_x and S_y of the spin Hamiltonian

return to the equilibrium tunnel mechanism



terms in S_x and S_y of the spin Hamiltonian

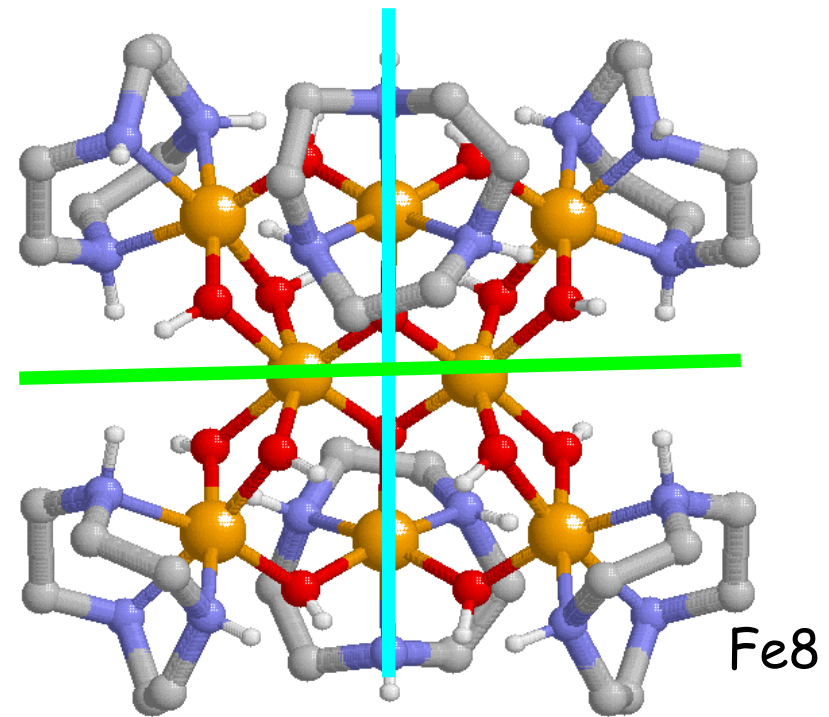
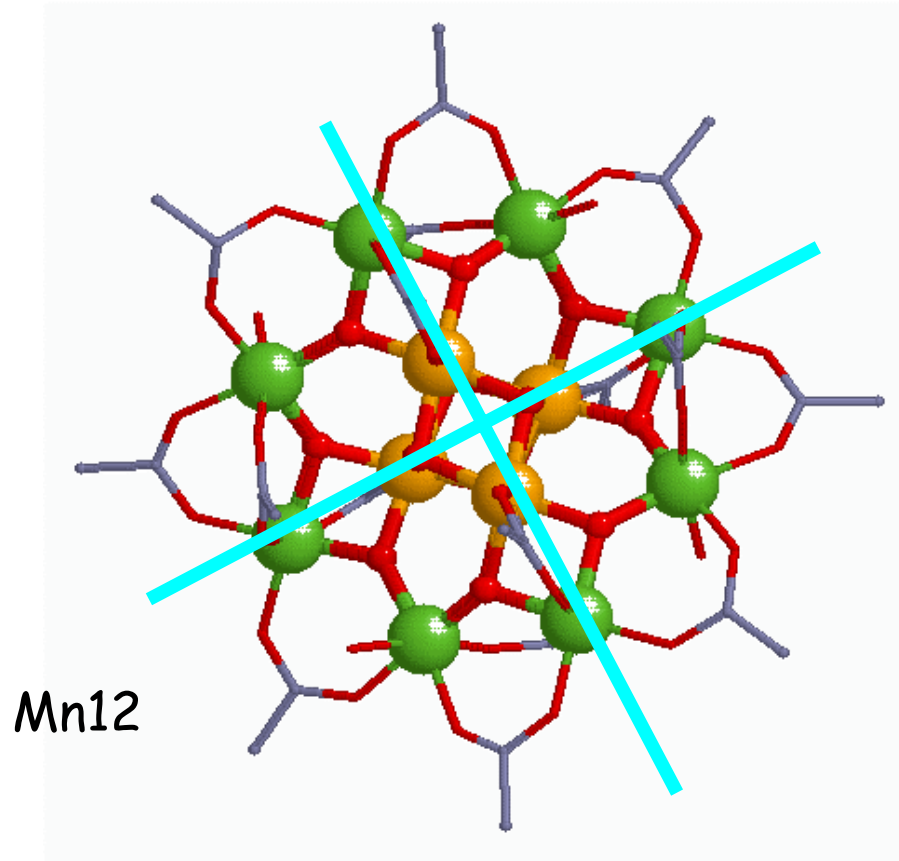
What is the difference ?



$$H = \mu_B \mathbf{S} \cdot \mathbf{g} \cdot \mathbf{B} - D S_z^2 + E (S_x^2 - S_y^2) + \cancel{BS_z^4} + C (S_+^4 + S_-^4)$$

Four fold axis
Tetragonal ($E=0$)

What is the difference ?



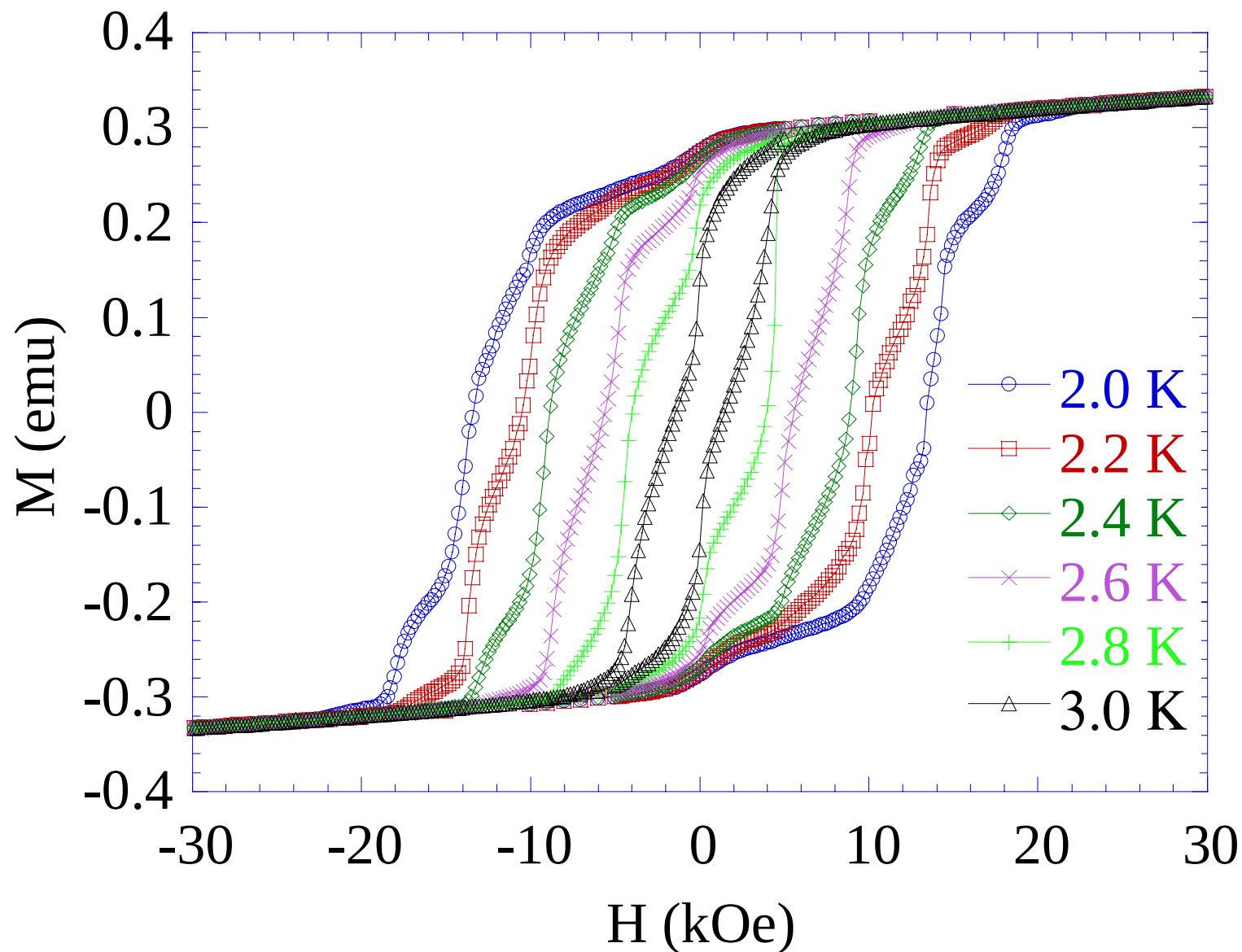
$$H = \mu_B \mathbf{S} \cdot \mathbf{g} \cdot \mathbf{B} - D S_z^2 + E (S_x^2 - S_y^2) + B S_z^4 + C (S_+^4 + S_-^4)$$

Four fold axis
Tetragonal ($E=0$)

Two fold axis
Rhombic ($E \neq 0$)

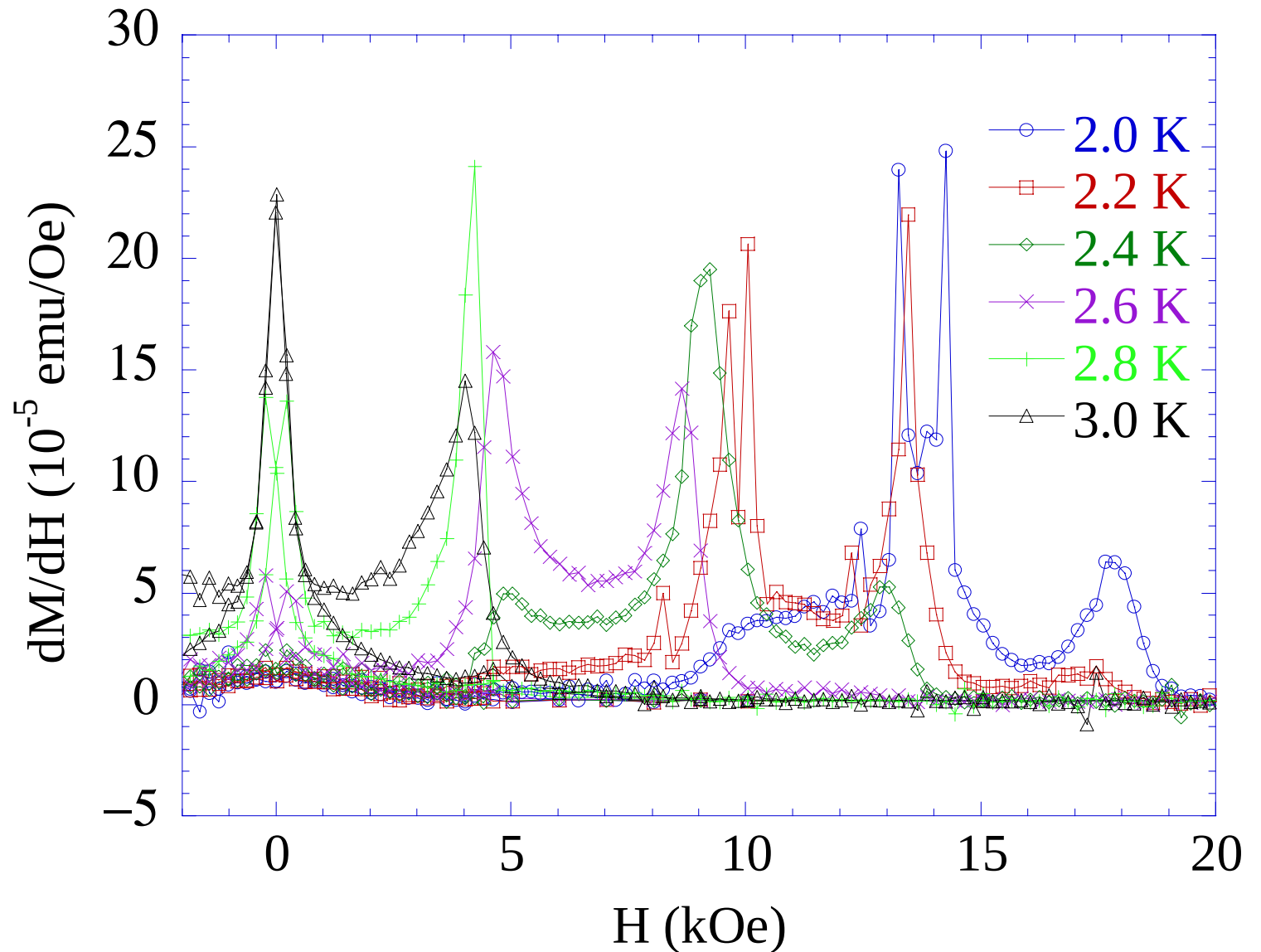
Hysteresis loops for Mn₁₂

Friedman et al.,
PRL, 1996;
Hernandez et al,
EPL, 1996;
Thomas et al.,
Nature, 1996

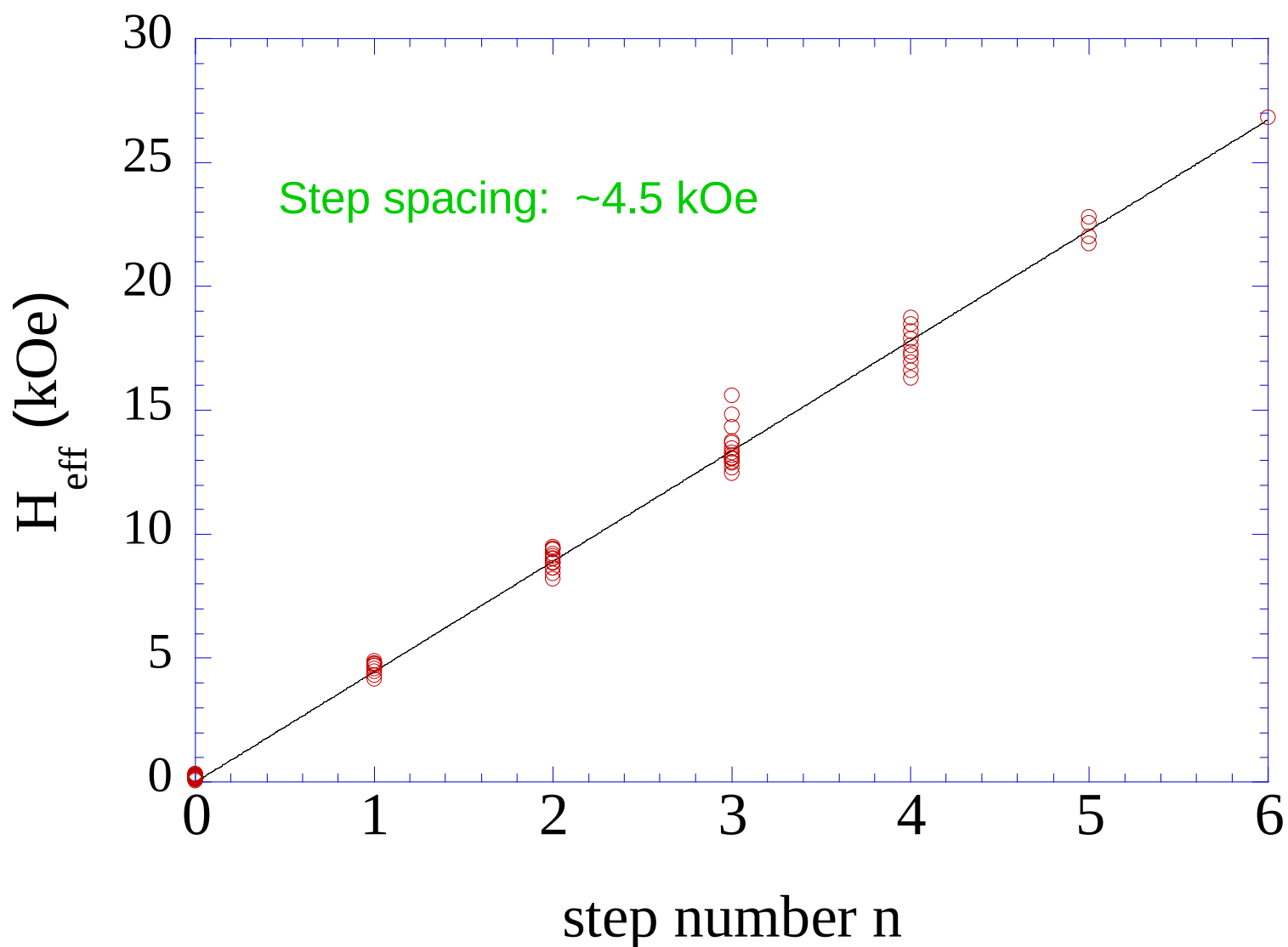


Hysteresis loops for Mn₁₂

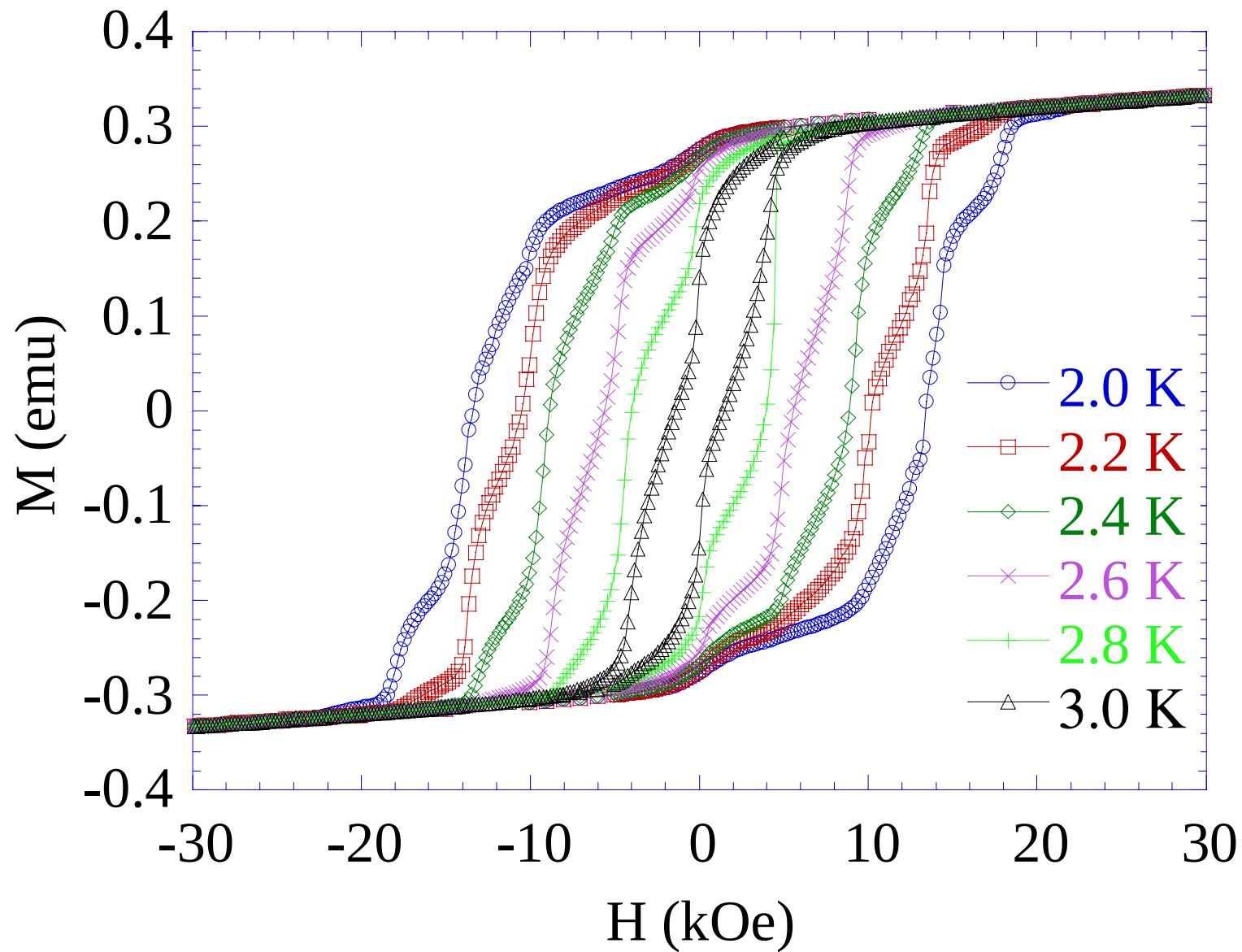
Friedman et al.,
PRL, 1996;
Hernandez et al.,
EPL, 1996;
Thomas et al.,
Nature, 1996



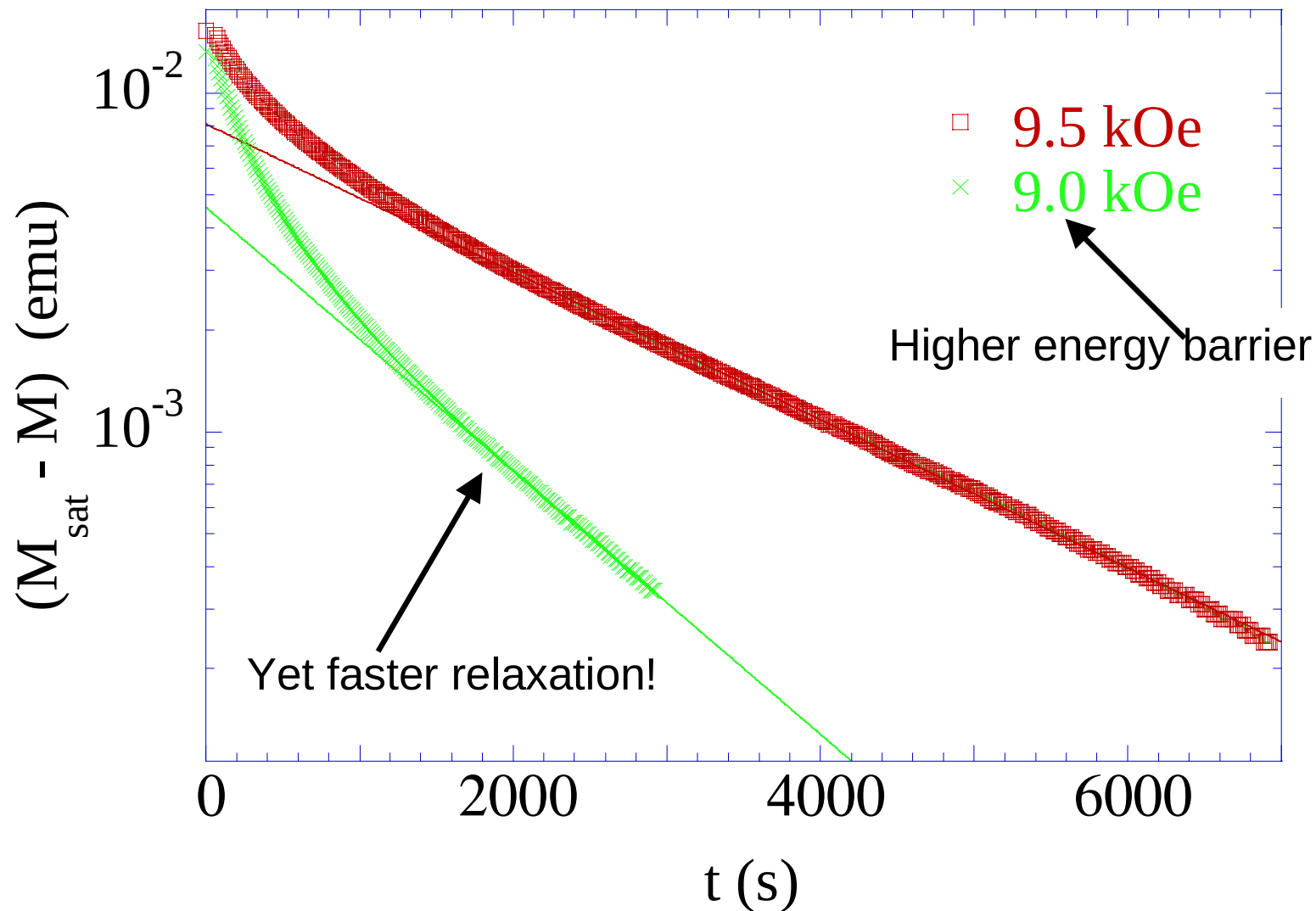
Uniform spacing between steps



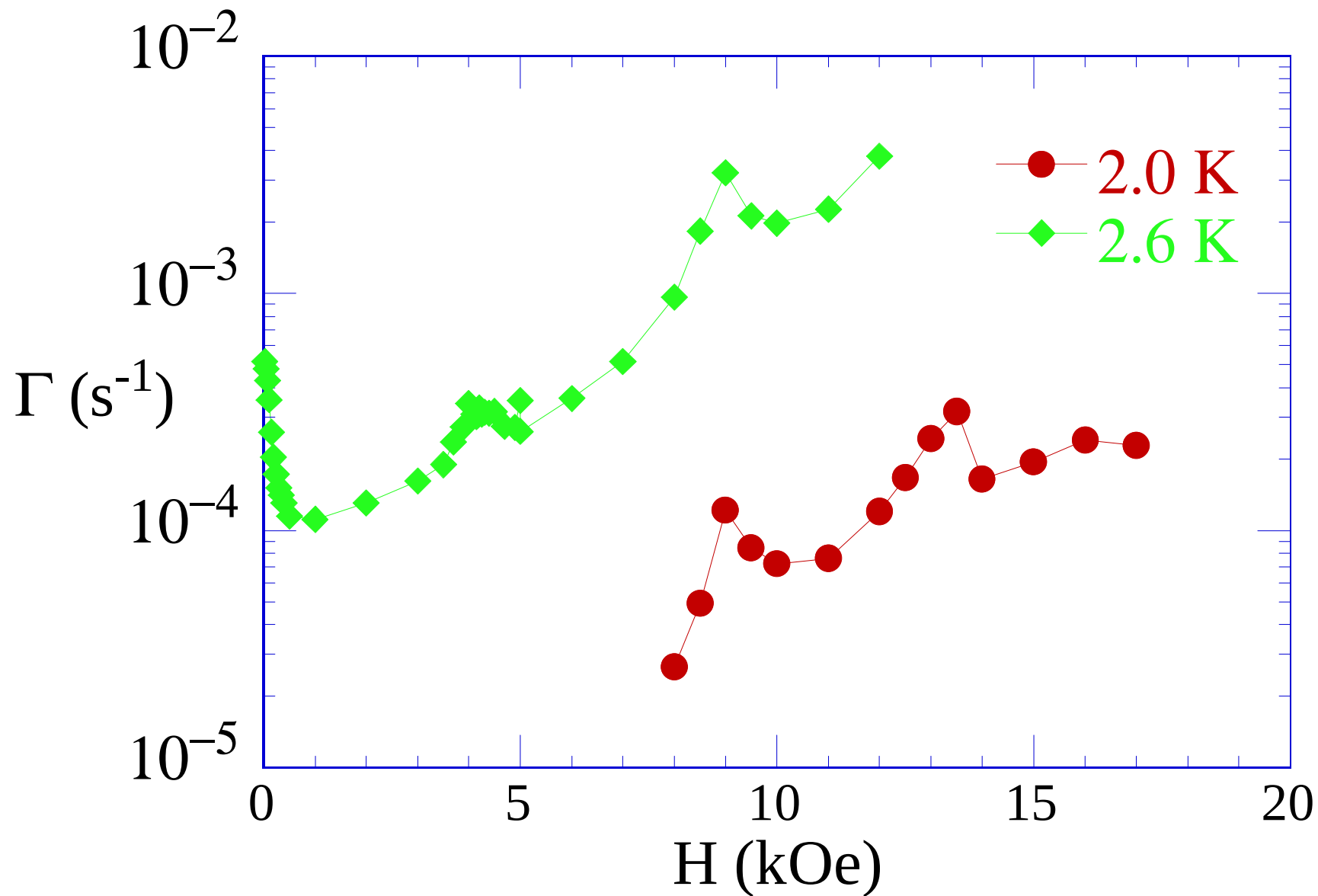
Hysteresis loops for Mn₁₂



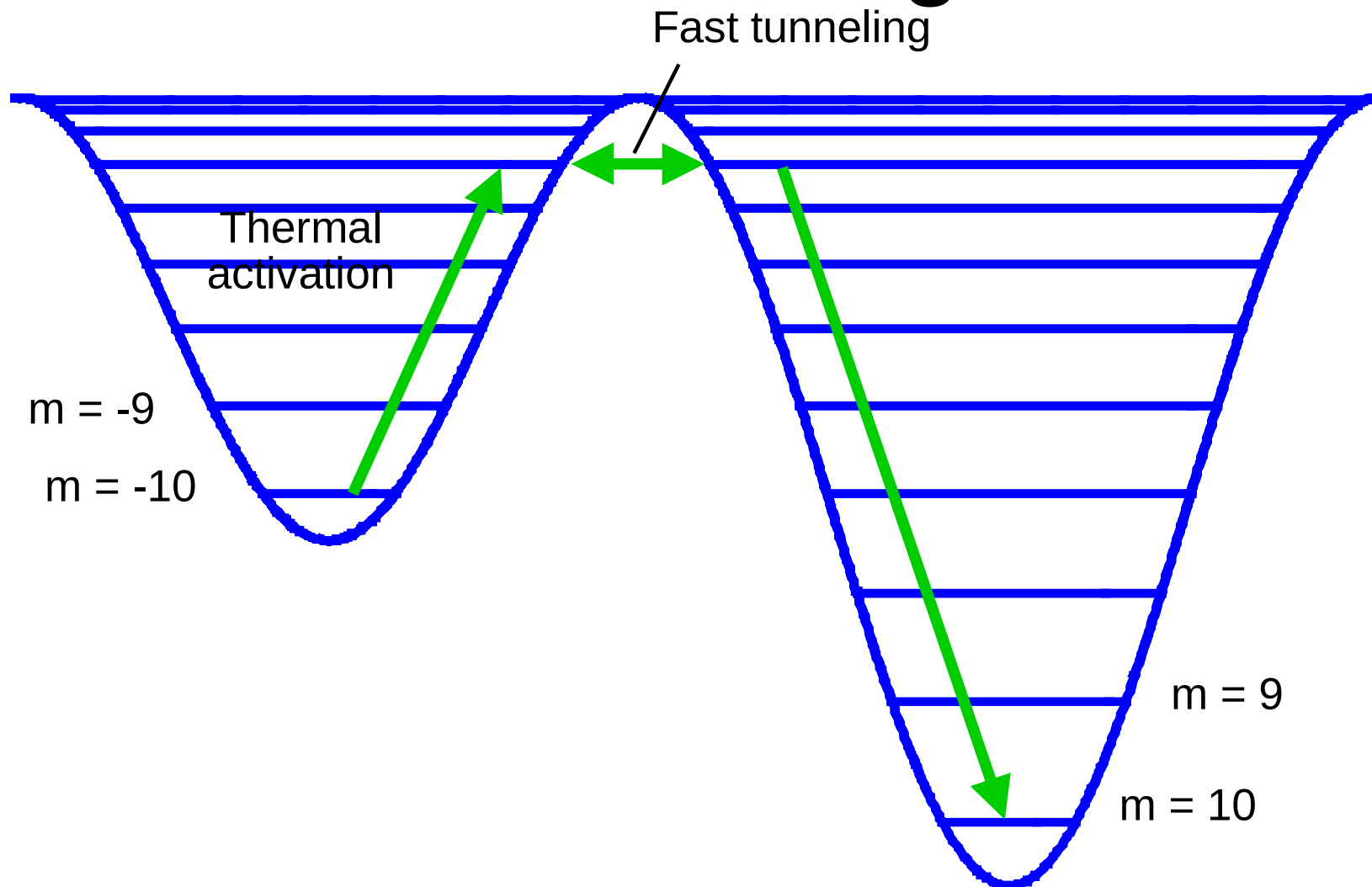
Enhanced Relaxation at Step Fields



Enhanced Relaxation at Step Fields



Thermally Assisted Resonant Tunneling



Tunneling occurs when levels in opposite wells align.

Hamiltonian for Mn₁₂

$$= -DS_z^2 - g\mu_B \mathbf{S} \cdot \mathbf{H}$$

The field at which $|m\rangle$ (in the left well) crosses $|-m+n\rangle$ (in the right well):

$$H_{m,-m+n} = \frac{-Dn}{g\mu_B}$$

Steps occur at regular intervals of field, as observed.

Step occurs every 4.5 kOe $\Rightarrow$ $D/g = 0.31$ K

Compare with ESR data:

D = 0.56 K, g = 1.93

(Barra et al., PRB, 1997)

D/g = 0.29 K

Hamiltonian for Mn₁₂

$$= -DS_z^2 - g\mu_B \mathbf{S} \cdot \mathbf{H} - BS_z^4$$

Spectroscopic studies revealed a 4th-order longitudinal anisotropy term $B \sim 1.1$ mK. (ESR: Barra et al., PRB, 1997 and Hill et al., PRL, 1998; INS: Mirebeau et al., PRL, 1999, Zhong et al., JAP, 2000 and Bao et al., cond-mat, 2000)

⇒ Different pairs of levels cross at slightly different fields.

⇒ Allows for the Examination of the Crossover from Thermally Assisted to Pure Quantum Tunneling.

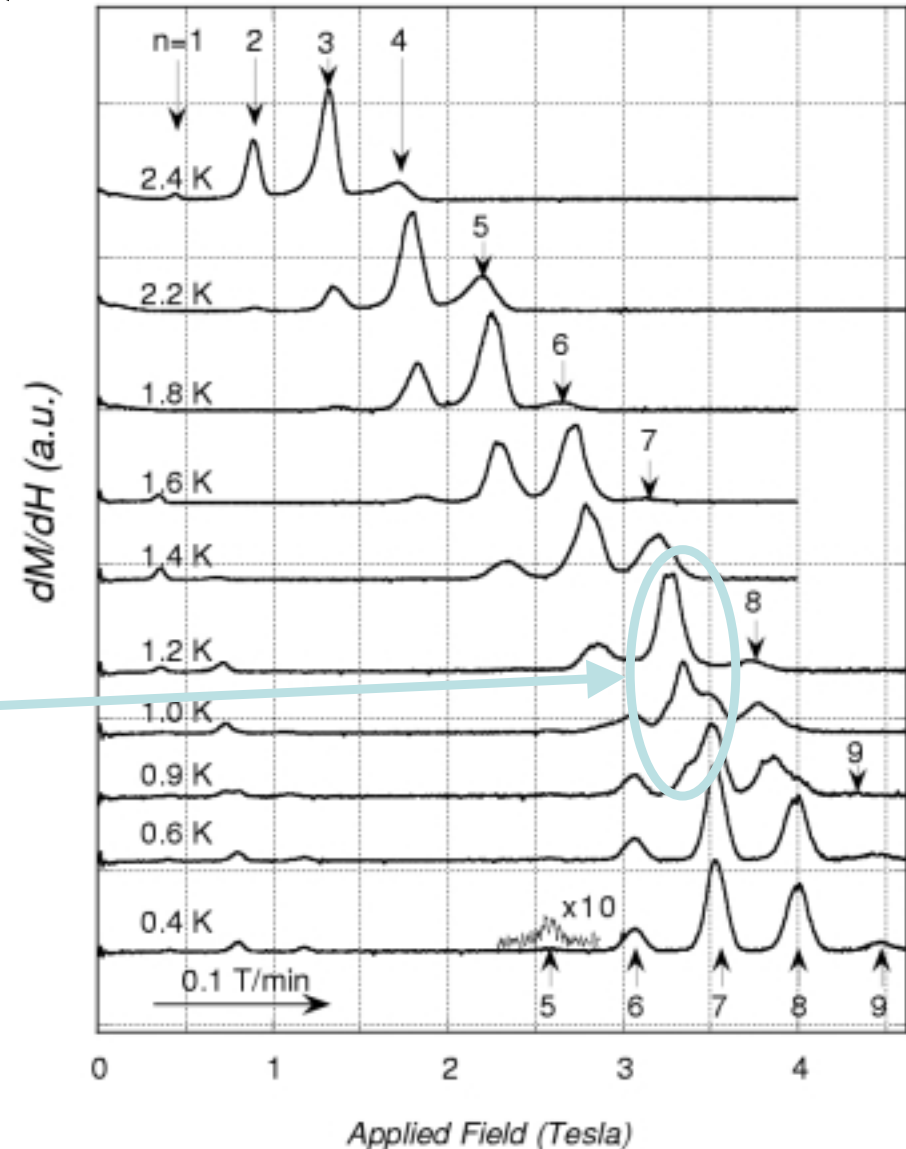
Crossover to Ground-state Tunneling

Level crossing fields:

$$H_{m,m'} = \frac{Dn}{g\mu_B} \left[1 + \frac{B}{D} (m^2 + m'^2) \right]$$

Abrupt “first-order” transition between thermally assisted and ground state tunneling.

Theory: Chudnovsky and Garanin, PRL, 1997; Exp't: Kent, et al., EPL, 2000, Mertes et al., JAP, 2001.

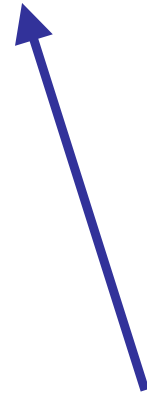


Fe₈ Hamiltonian in Zero Field

$$= -DS_z^2 + E(S_x^2 - S_y^2) + C(S_+^4 + S_-^4)$$



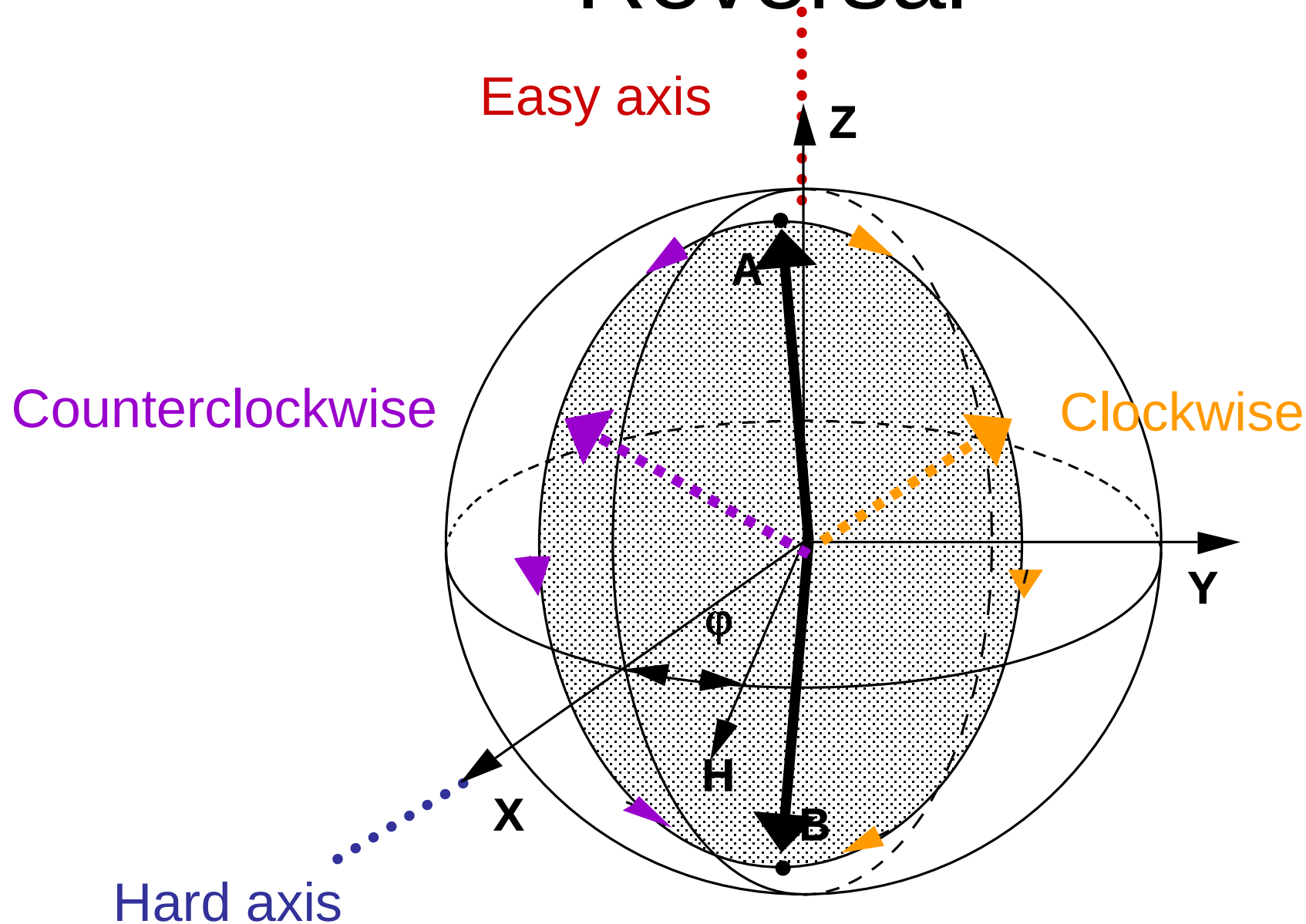
Easy Axis



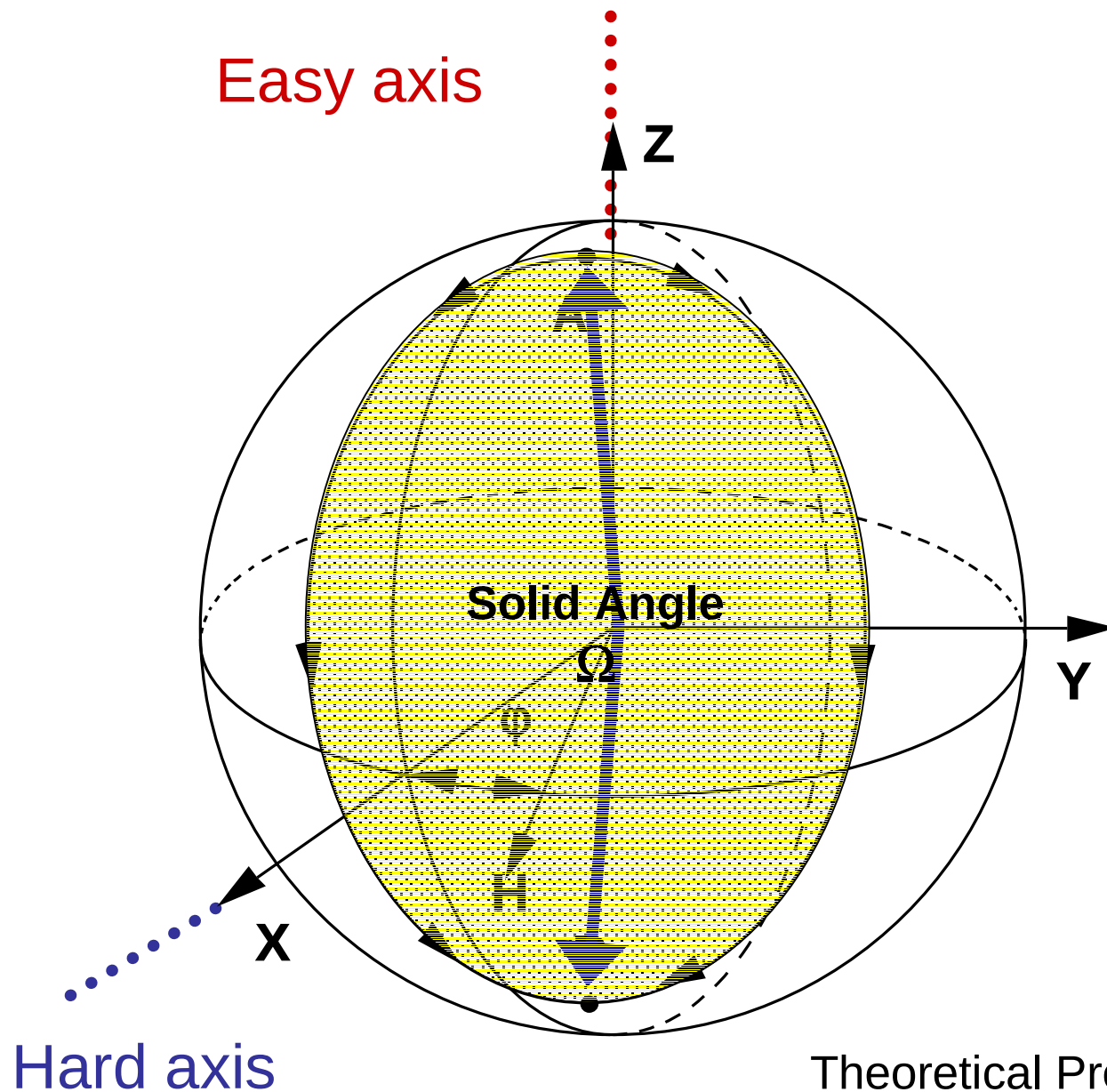
Hard Axis

Spin wants to rotate in the y-z plane

Two Paths for Magnetization Reversal



Destructive Topological Interference



Equivalence between paths is maintained when **H** is applied along the Hard Axis.

Topological (Berry's) phase depends on solid angle Ω encribed by the two paths.

Complete destructive interference occurs for certain discrete values of Ω .

Theoretical Prediction: A. Garg., 1993.

Destructive Topological Interference

Modulation of Tunnel Splitting:

$$\Delta = \cos(S\Omega),$$

where Ω depends on the field along the Hard Axis.

When $S\Omega = \pi/2, 3\pi/2, 5\pi/2\dots$, tunneling is completely suppressed!

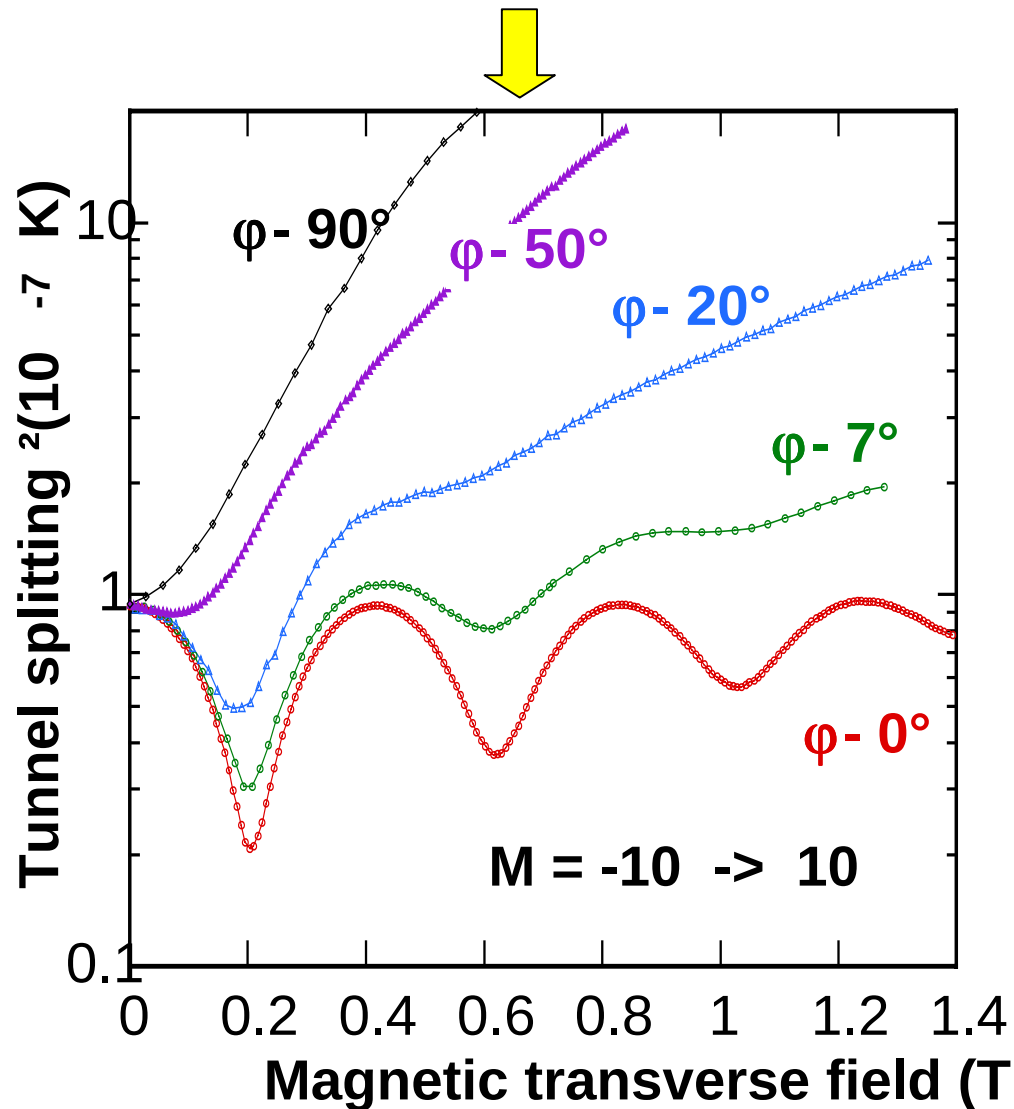
Interval between such destructive interference points:

$$\Delta H = \frac{2}{g\mu_B} \sqrt{2E(E + D)}$$

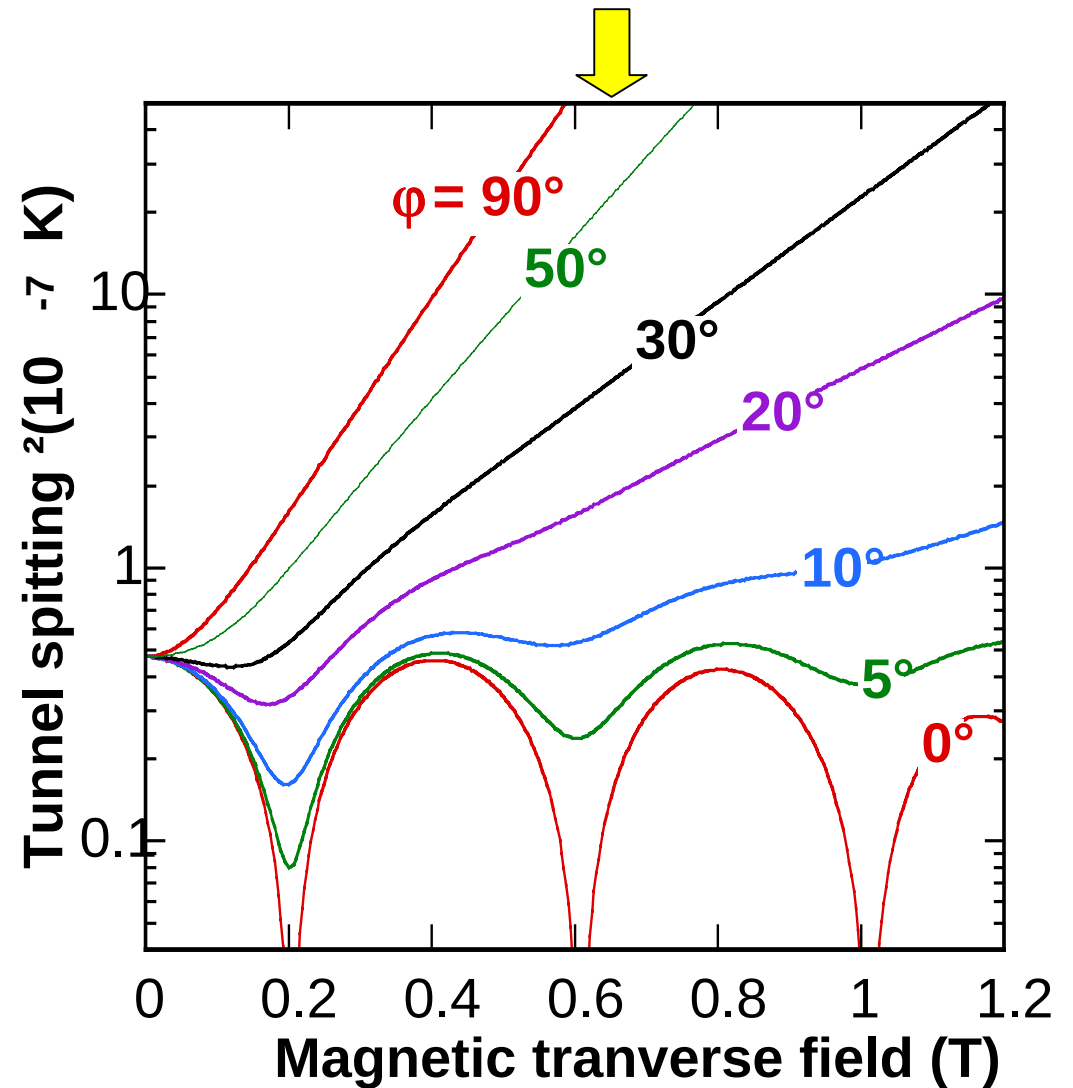
A. Garg., 1993.

Measured Tunnel Splitting

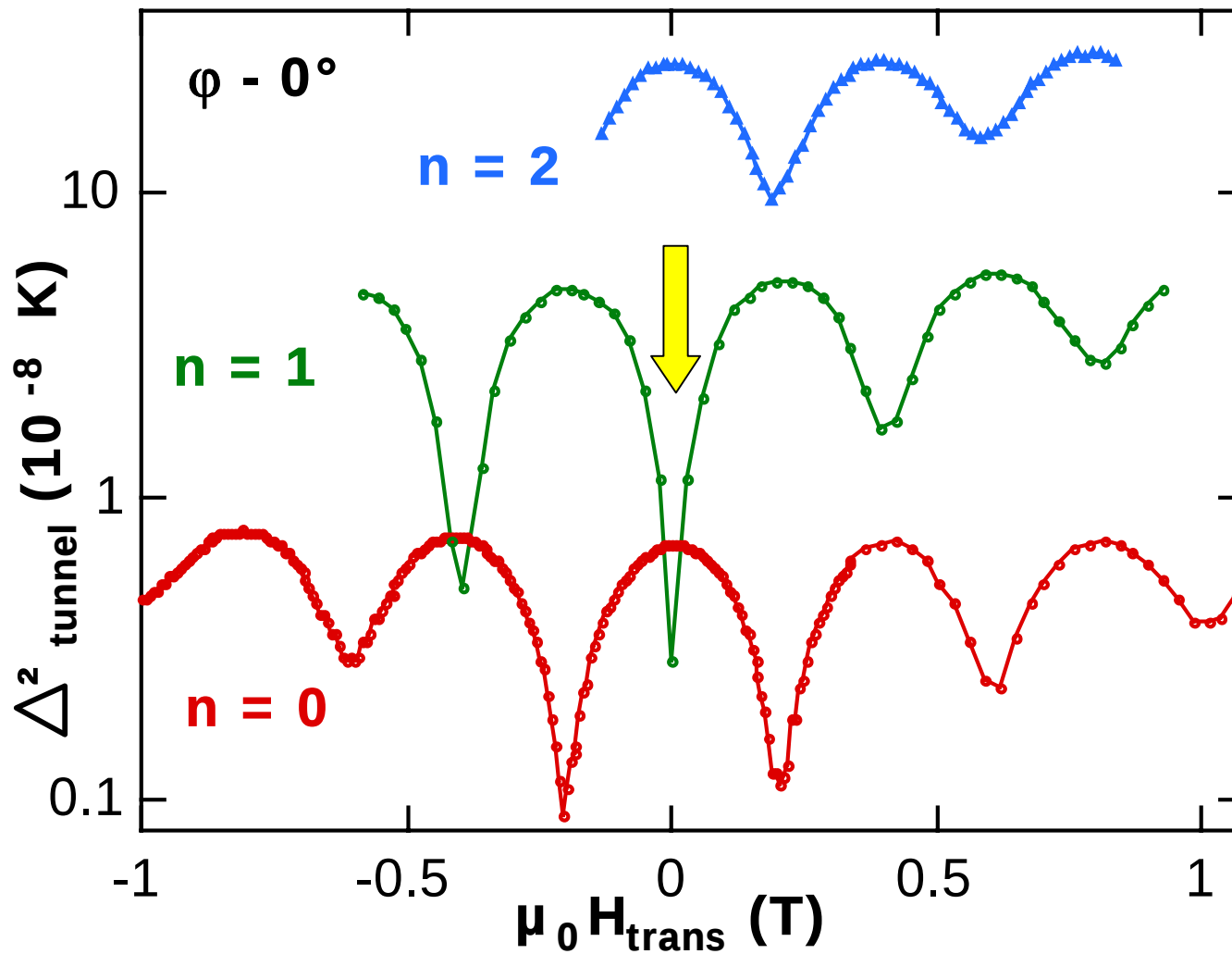
experimental



calculated with
 $D = -0.29, E = 0.046, C = -2.9 \times 10^{-5} \text{ K}$



Parity Effect: Odd vs. Even Resonances



W. Wernsdorfer and R. Sessoli, Science, 1999.

What Causes Tunneling and Why the Parity Effect in Fe₈

- Tunneling is produced by terms in the Hamiltonian that do not commute with S_z .

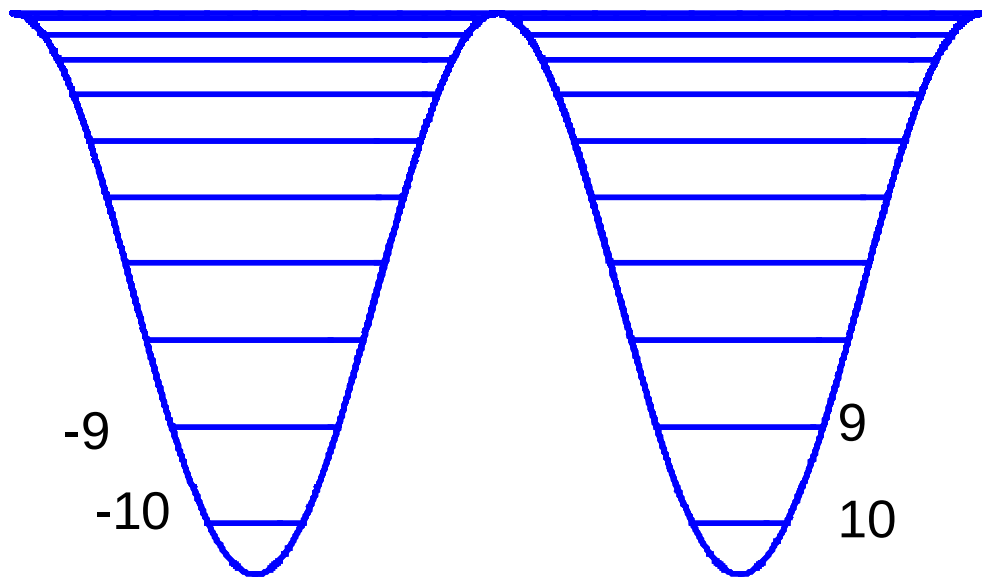
- For Fe₈, these terms are

$$E(S_x^2 - S_y^2) = \frac{E}{2}(S_+^2 + S_-^2)$$

- Selection rule: $\Delta m = \pm 2p$ ($p = 1, 2, 3, \dots$)
- Every other tunneling resonance is forbidden!

What Causes Tunneling and Why the Parity Effect in Fe₈

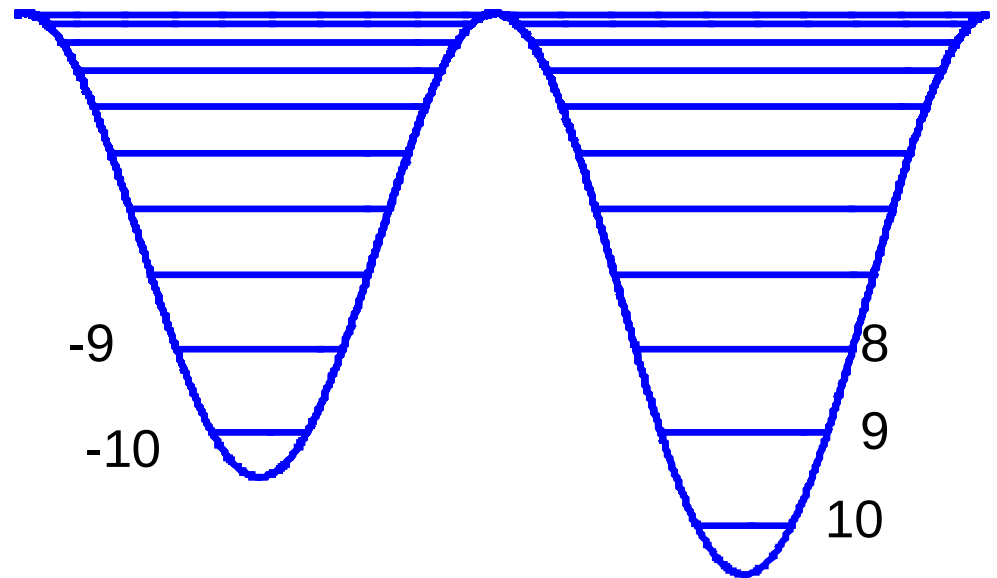
$n = 0$



$$\Delta m = \pm 2p$$

Tunneling Allowed

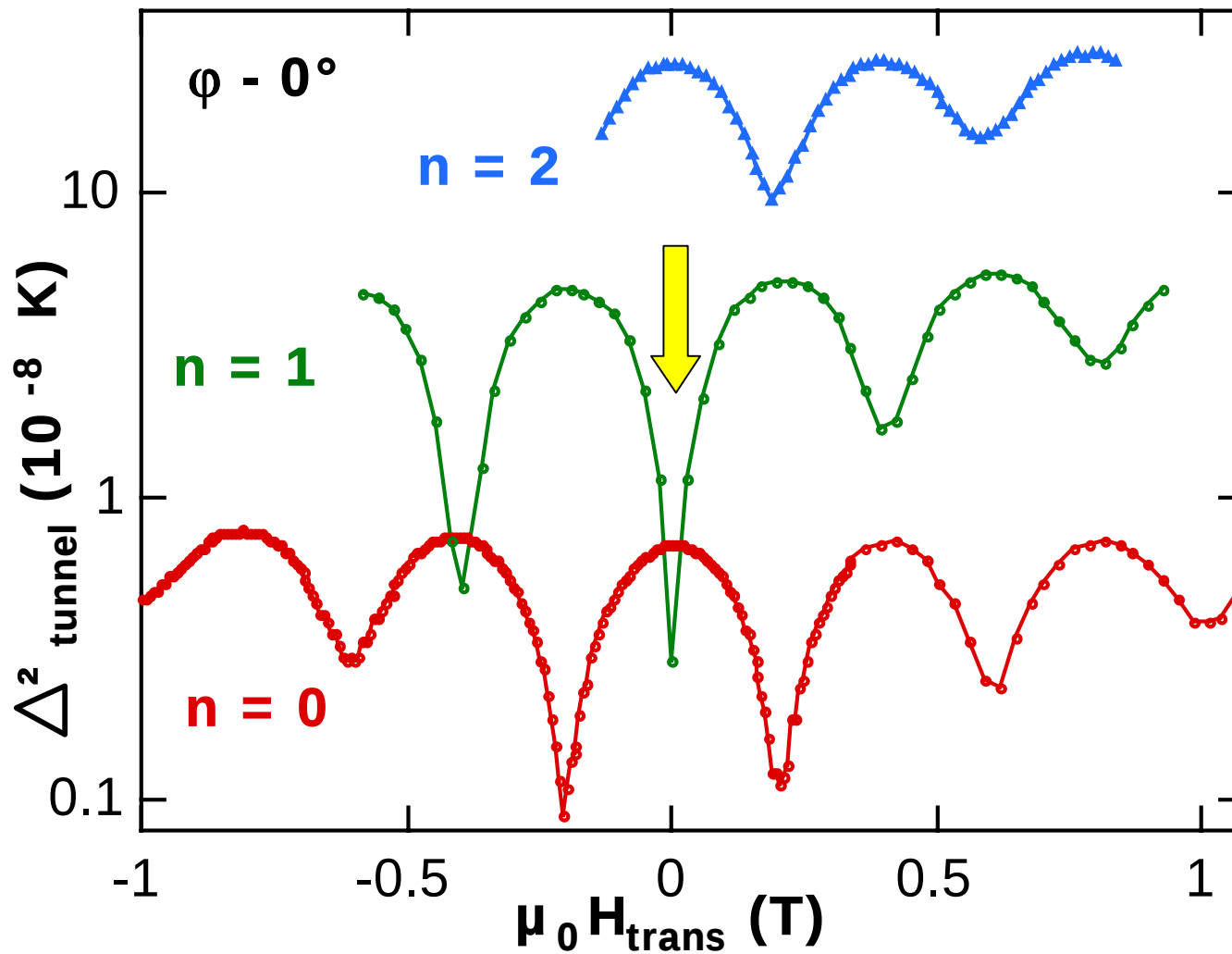
$n = 1$



$$\Delta m \neq \pm 2p$$

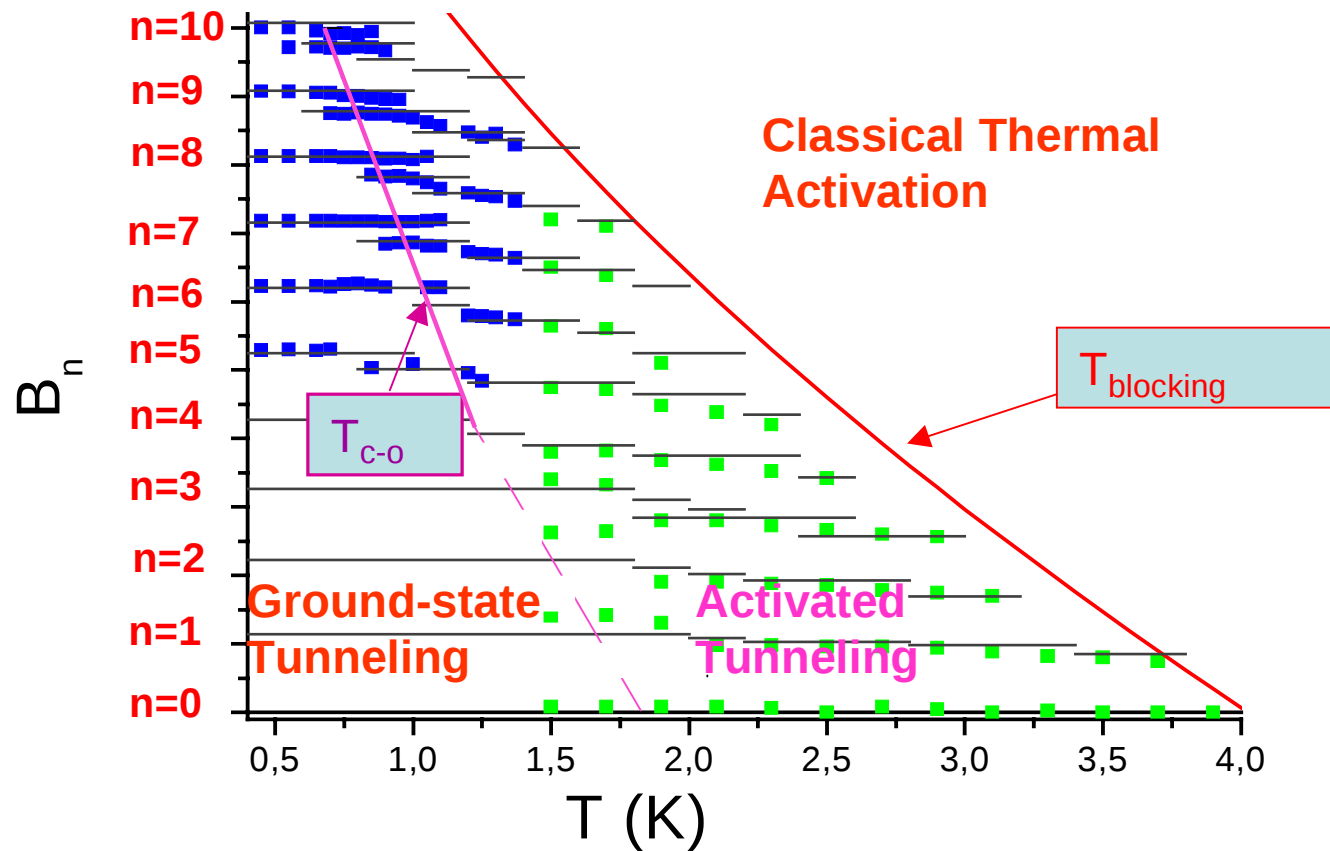
Tunneling Forbidden

Parity Effect: Odd vs. Even Resonances



W. Wernsdorfer and R. Sessoli, Science, 1999.

Crossover From Classical to Quantum Regime (Mn₁₂-ac)



Measured (■ ■) and Calculated (—) Resonance Fields

Barbara et al, JMMM 140-144, 1891 (1995) and J. Phys. Jpn. 69, 383 (2000)

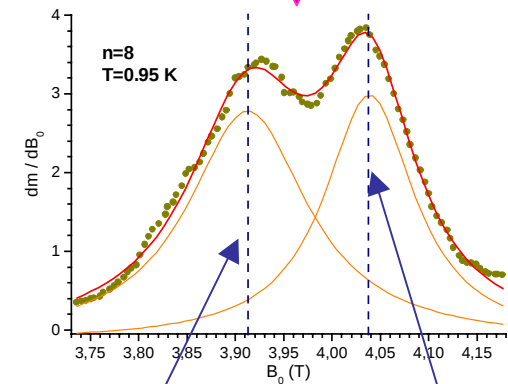
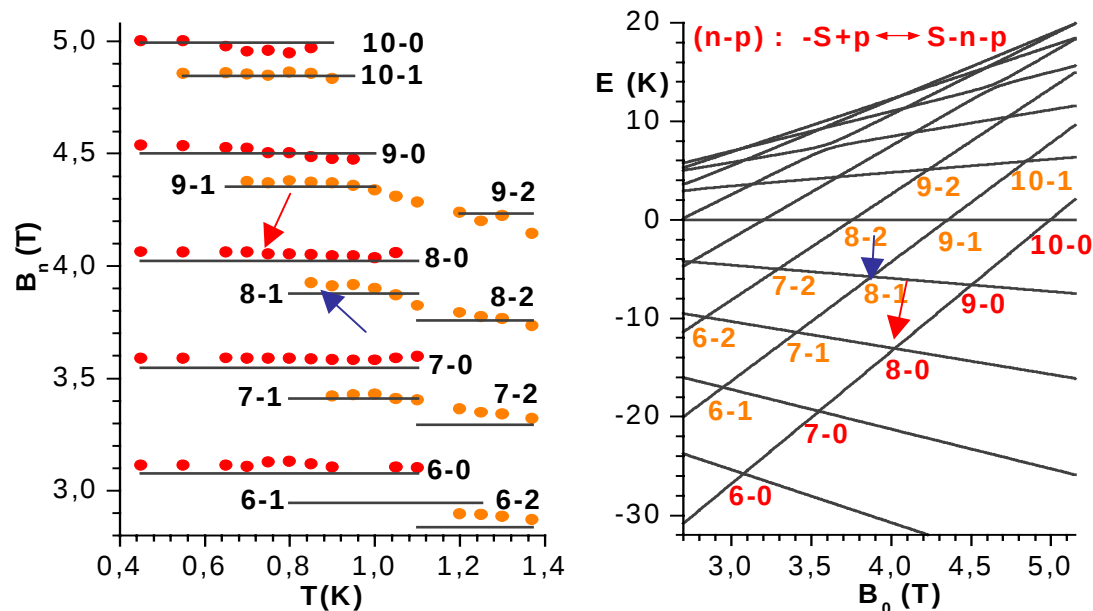
Paulsen, et al, JMMM 140-144, 379 (1995); NATO, Appl. Sci. 301, Kluwer (1995)

The Tunnel Window: An effect of weak Hyperfine Interactions

Data points and calculated lines

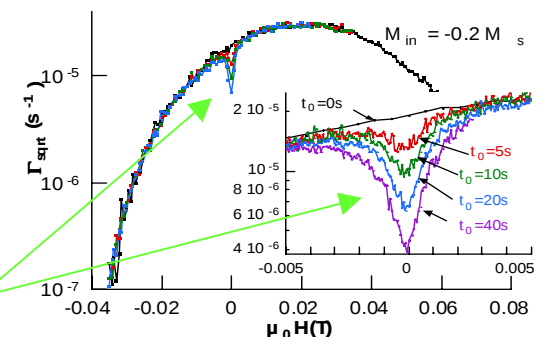
Level Scheme

Inhomogeneous broadening of
Two resonances: Dipolar fields



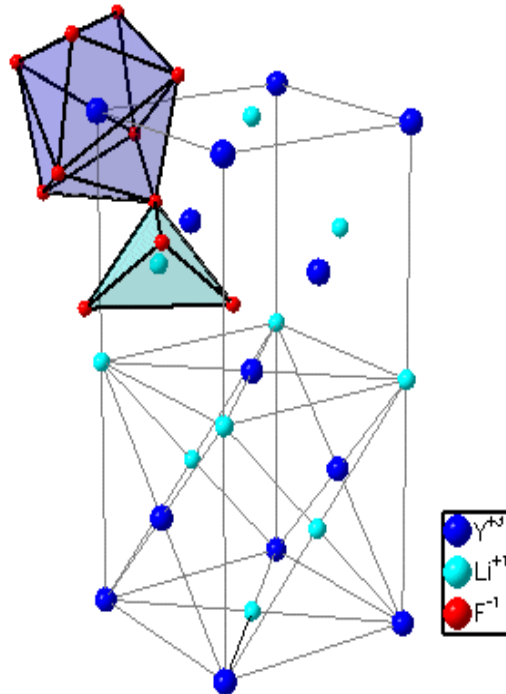
- Chiorescu et al, PRL, 83, 947 (1999)
- Barbara et al, J. Phys. Jpn. 69, 383 (2000)
- Kent et al, EPL, 49, 521 (2000)
- Wernsdorfer et al, PRL (1999)

Homogeneous
broadening of nuclear
spins: Tunnel window



Effects of Strong Hyperfine Interactions:

Case of Rare-earth ions: Ho^{3+} in $\text{Y}_{0.998}\text{Ho}_{0.002}\text{LiF}_4$



Tetragonal symmetry
(Ho in S_4)

$$J = L + S = 8; \quad g_J = 5/4$$

Dipolar interactions between $\text{Ho}^{3+} \rightarrow \ll \text{mT}$

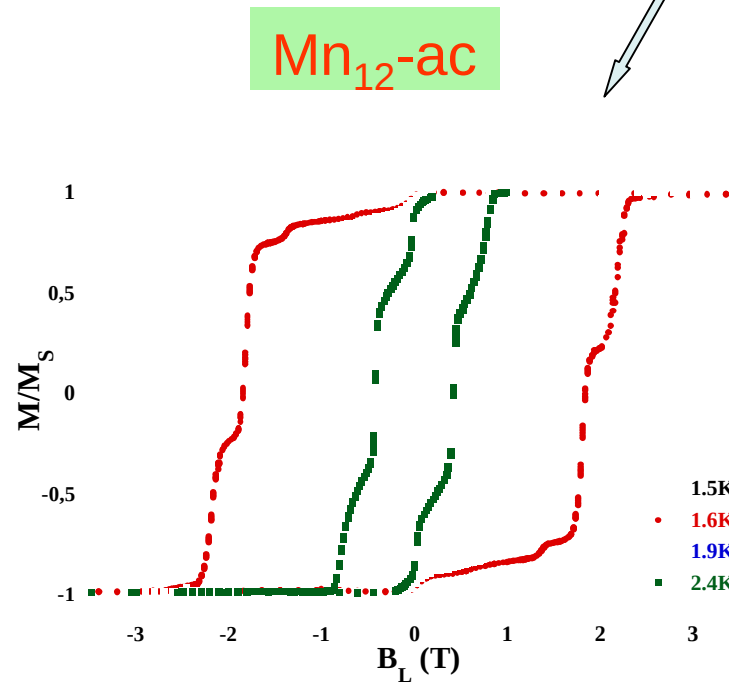
$$H_{\text{CF-Z}} = -B_2^0 O_2^0 - B_4^0 O_4^0 - B_4^4 O_4^4 - B_6^0 O_6^0 - B_6^4 O_6^4 - g_J \mu_B \mathbf{JH}$$

B_l^m : accurately determined by high resolution optical spectroscopy

Sh. Gifeisman et al, Opt. Spect. (USSR) 44, 68 (1978); N.I. Agladze et al, PRL, 66, 477 (1991)

Hysteresis loop of Ho^{3+} ions in YLiF_4

Comparison with $\text{Mn}_{12}\text{-ac}$

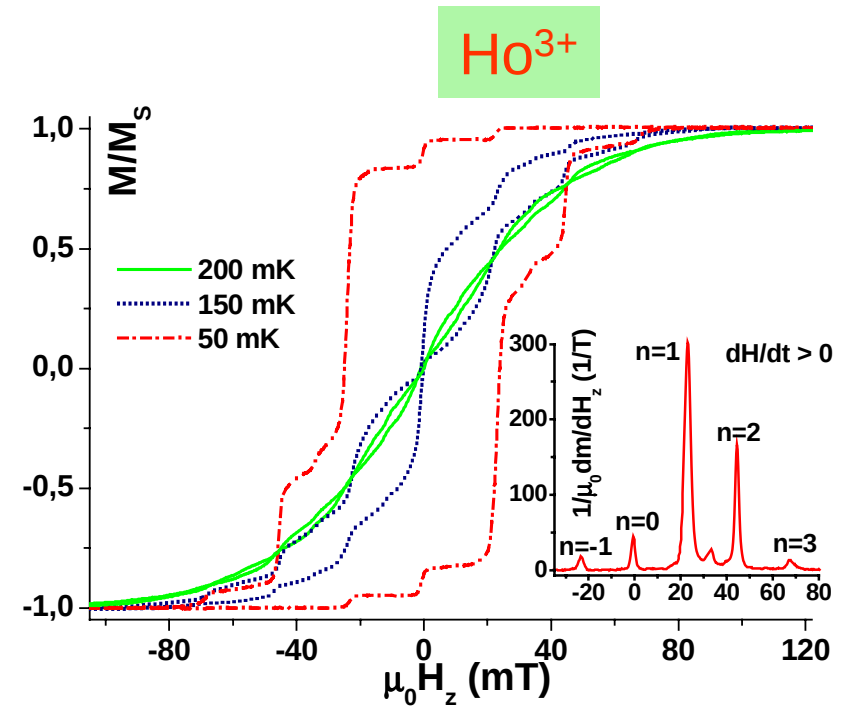


Thomas et al, Nature (1996)

Friedman et al, PRL (1996), Hernandez et al, EPL (1996)

Steps at $B_n = 450.n$ (mT)

Tunneling of $\text{Mn}_{12}\text{-ac}$ Molecules



Giraud et al, PRL, 87, 057203-1 (2001)

Steps at $B_n = 23.n$ (mT)

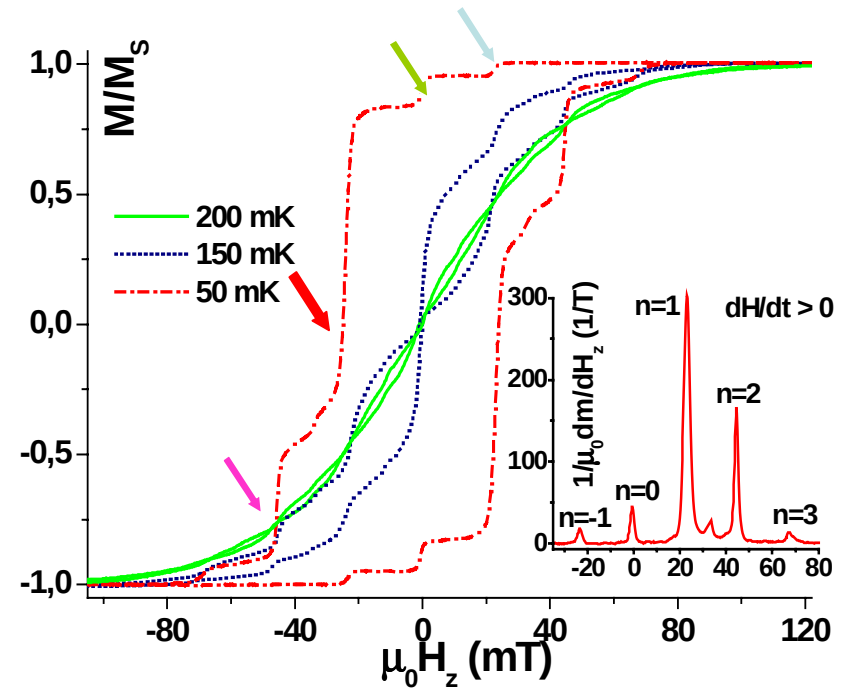
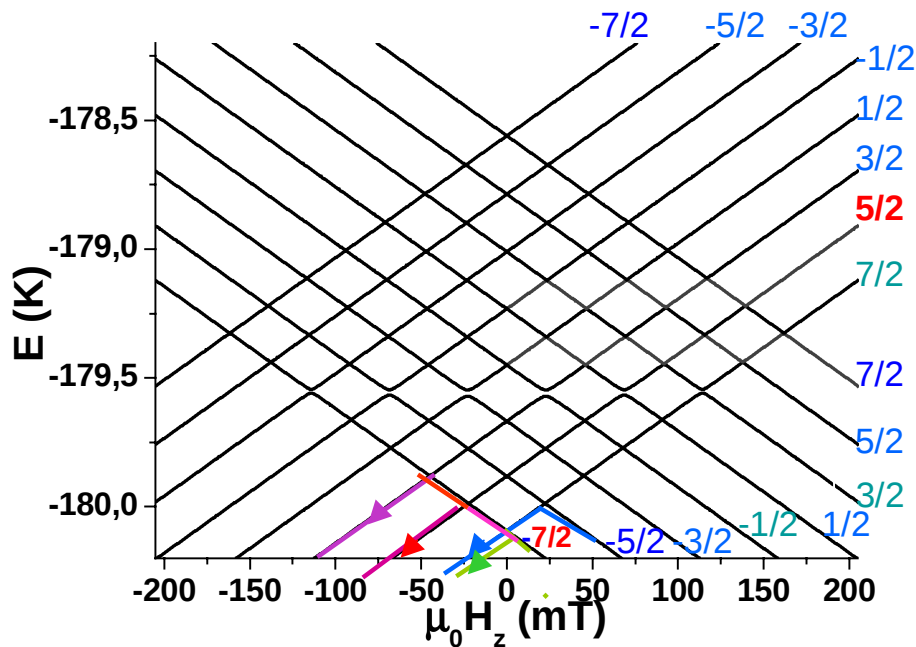
Tunneling of Ho^{3+} ion

Role of Strong Hyperfine Interactions

$$H = H_{\text{CF-Z}} + A.I.J$$



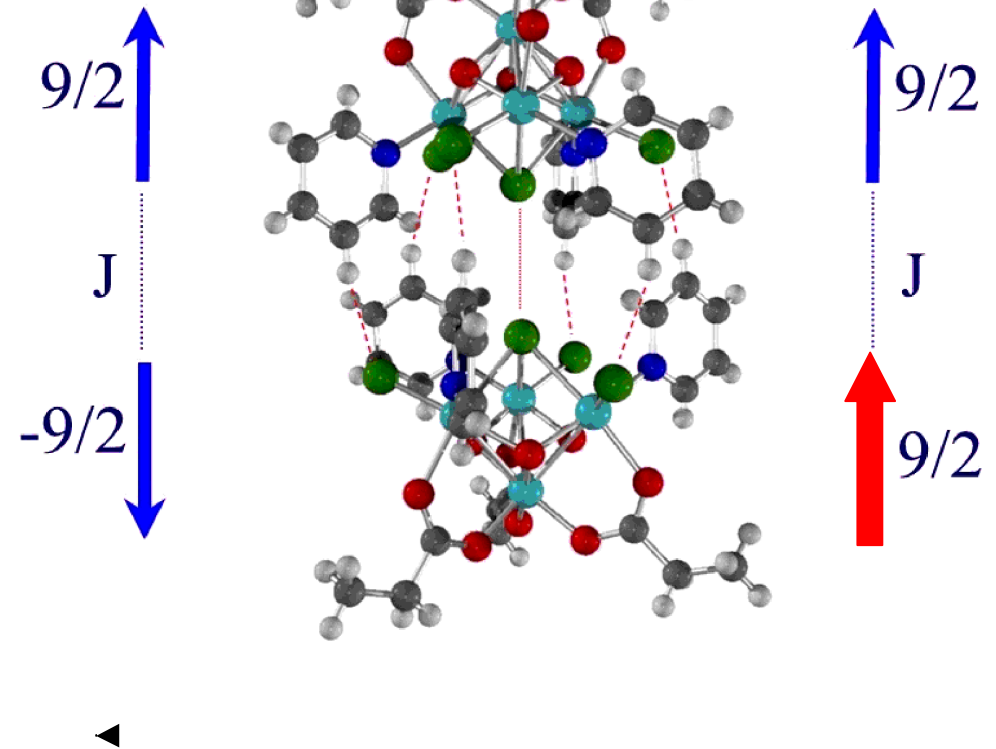
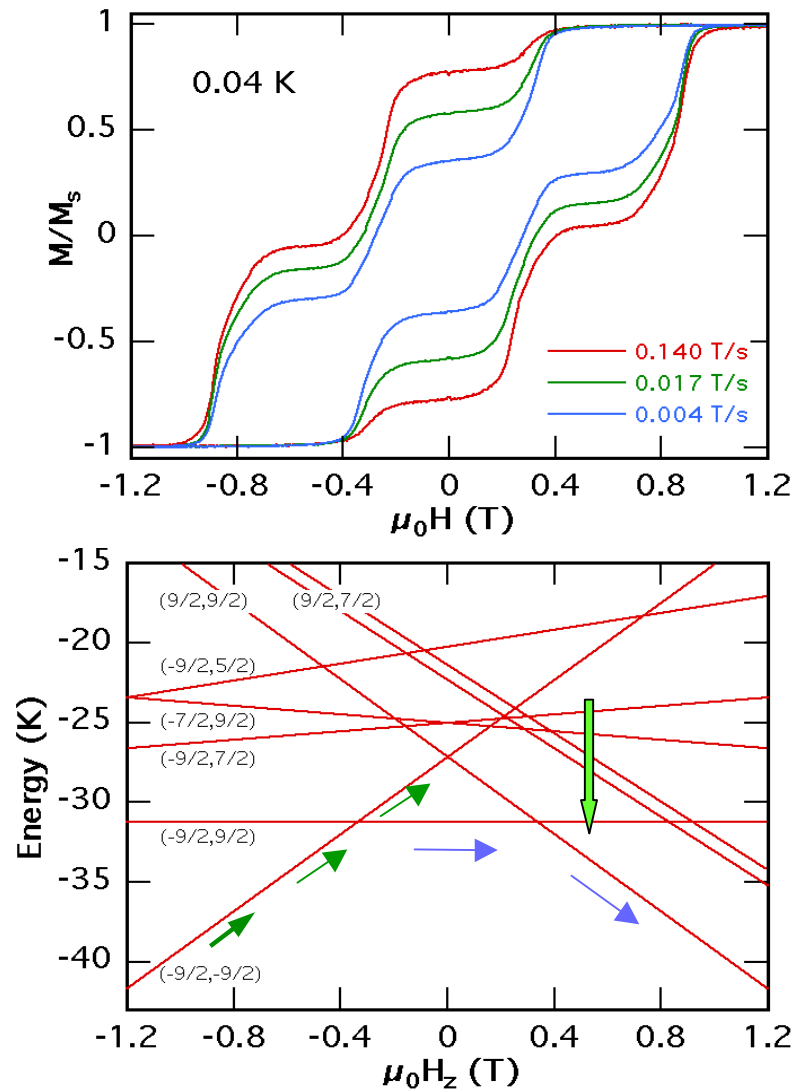
Induce Tunneling of Electronic Moments



Co-Tunneling of Electronic and Nuclear Spins:
Electro-nuclear entanglement

Avoided Level Crossings between $|\Psi-, I_z\rangle$ and $|\Psi+, I_z'\rangle$ if $|I_z - I_z'|/2$ integer

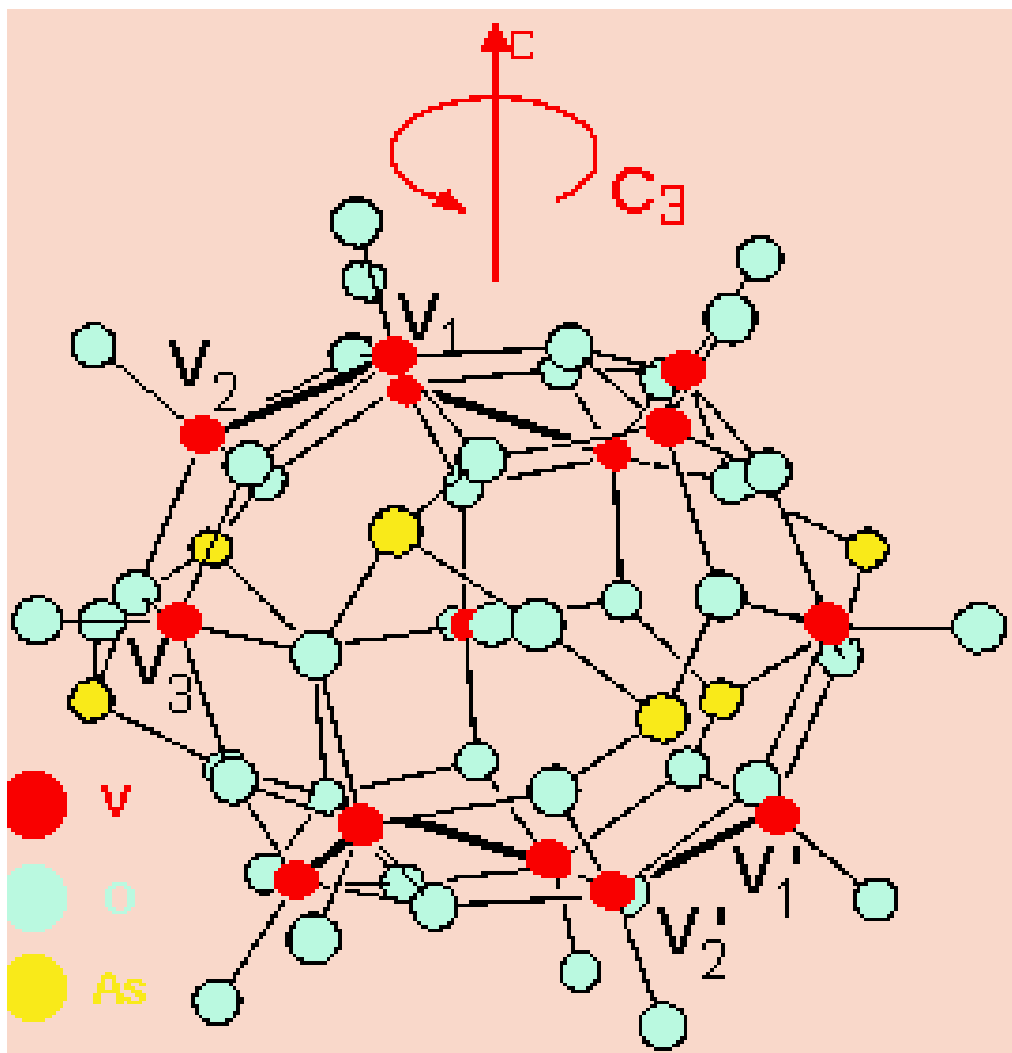
Exchange-biased quantum tunnelling in a dimer of Mn_4 molecule



W. Wernsdorfer et al, *Nature* **416**, 406 (2002)

V_{15} : The Archetype of Low spin Molecules

A Mesoscopic Spin $S=1/2$



Exchange interactions:
Antiferromagnetic $\sim$ several $10^2 K$

Müller, Döring, Angew. Chem. Intl.
Engl., 27, 171 (1988)

Anisotropy of g-factor:
 $\sim 0.6\%$

Ajiro et al, J..Low. Temp.
Phys. to appear (2003)

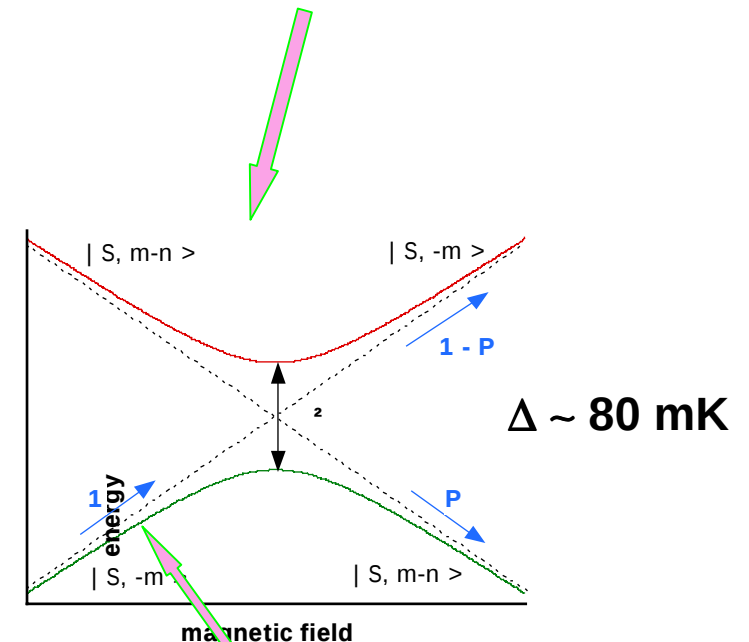
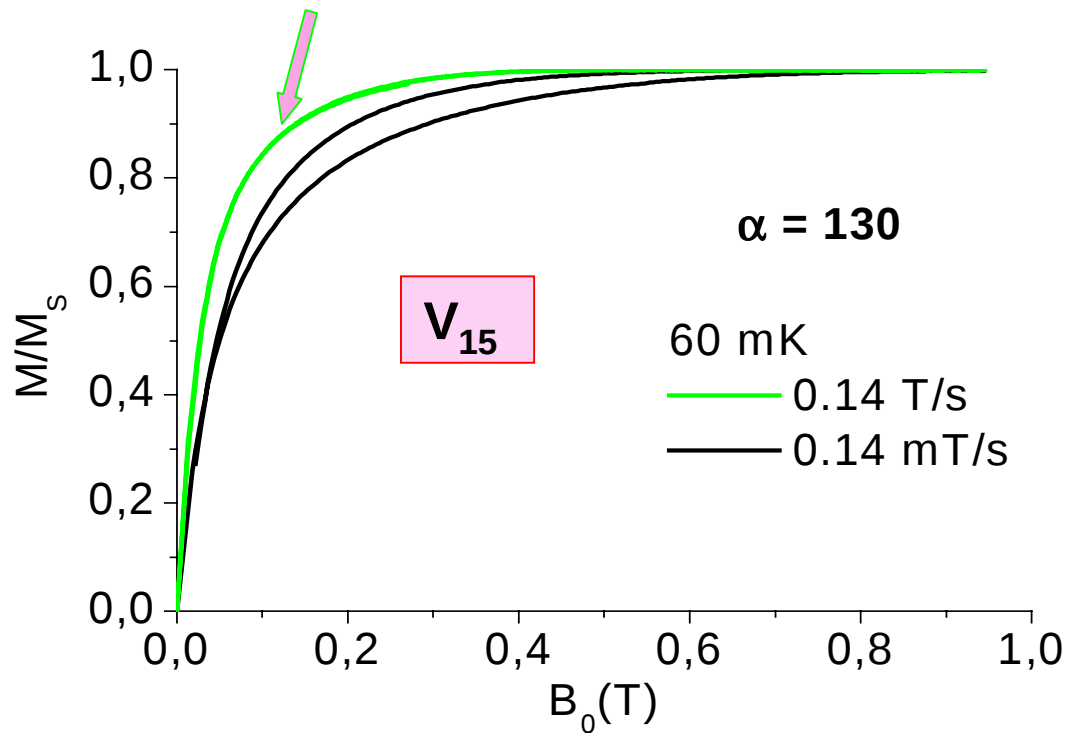
Barra et al, J. Am. Chem.
Soc. 114, 8509 (1992)

« Isolated V_{15} » :
A two-level system « without dissipation »

Fast sweeping rate / Weak coupling to the cryostat

Adiabatic Landau-Zener Spin Rotation

$M(H)$: Reversible and out of equilibrium



$$M(H) = dE(H)/dH$$

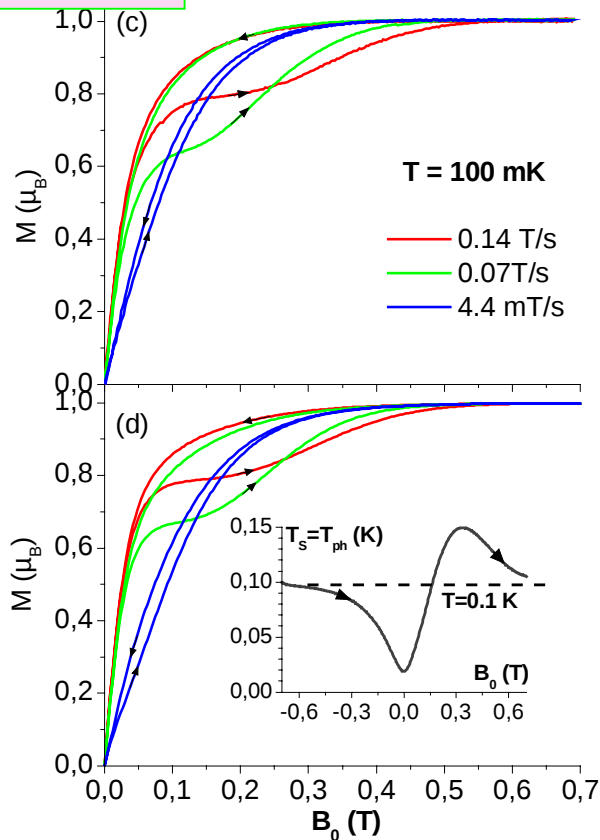
Nuclear Spin-Bath :
Weak Level Broadening

« Non-Isolated V_{15} » : A two-level system « with dissipation »

Low sweeping rate / Strong coupling to the cryostat

M(H): Irreversible

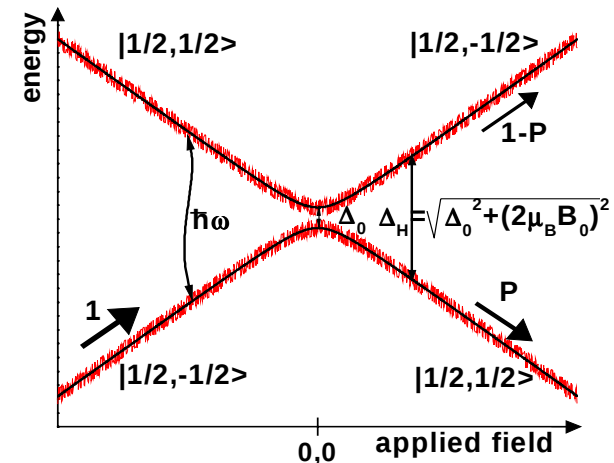
Measured



Calculated



LZS transition at Finite Temperature
(dissipative)



Phonon-bath → bottleneck model

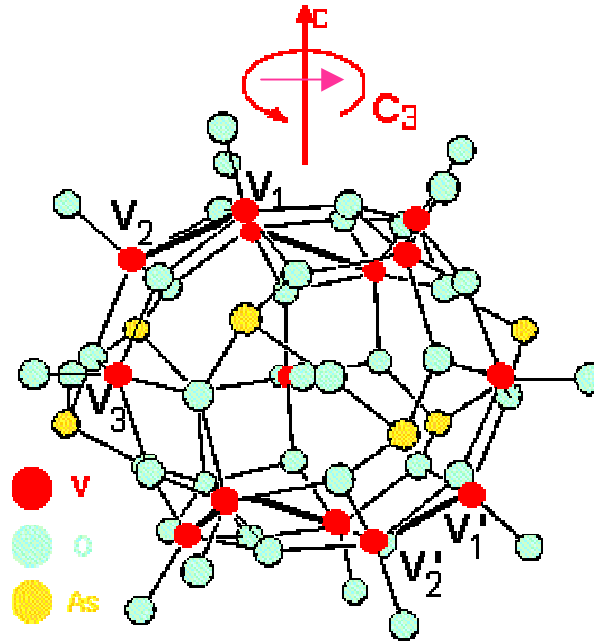
Abragam, Bleaney, 1970; Chiorescu et al, 1999.

Nuclear spin-bath → level broadening

Stamp, Prokofiev, 1998.

Chiorescu et al, PRL 84, 3454 (2000)

V_{15} : a Gapped Spin $\frac{1}{2}$ Molecule



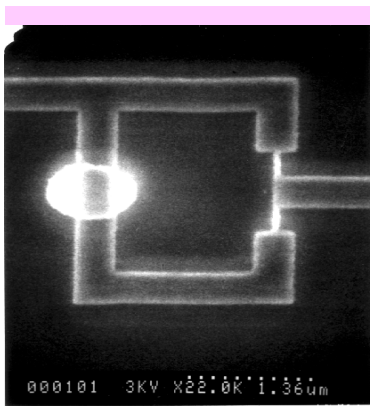
Time Reversal Symmetry $\rightarrow D = 0$ (Kramers Theorem)

The Multi-Spin Character of the Molecule
(15 spins)

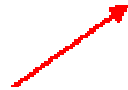
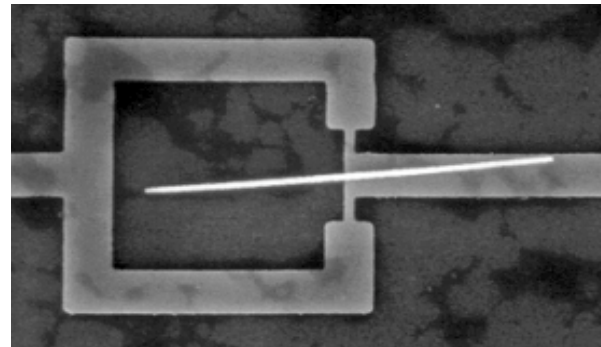
+

Dzyaloshinsky-Moriya interactions: $H_{DM} = -\sum_{ij} \mathbf{D}_{ij} \mathbf{S}_i \times \mathbf{S}_j$

Magnetism: From Macroscopic to Single atoms

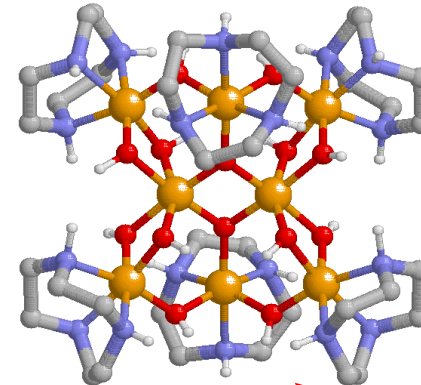


submicron



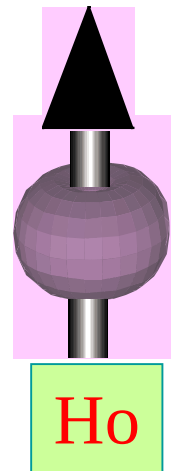
Nano-wires

Nano-particles

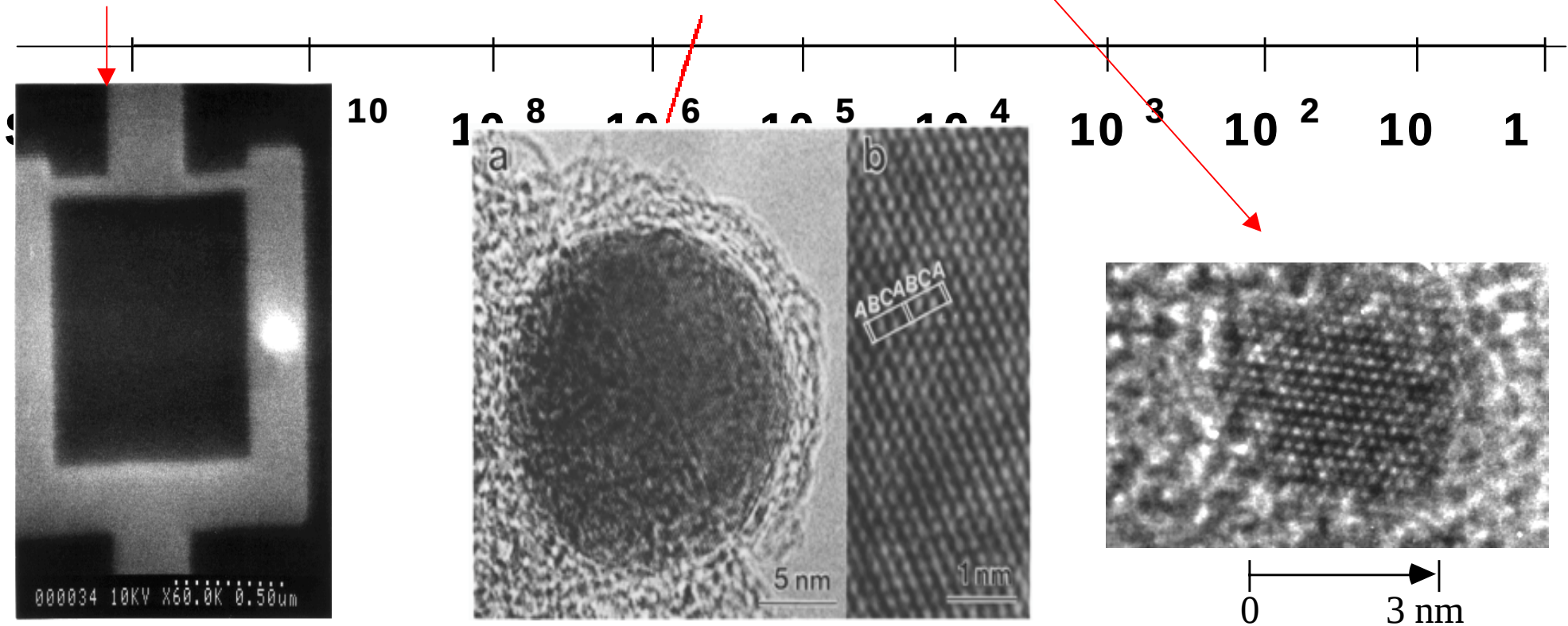


clusters

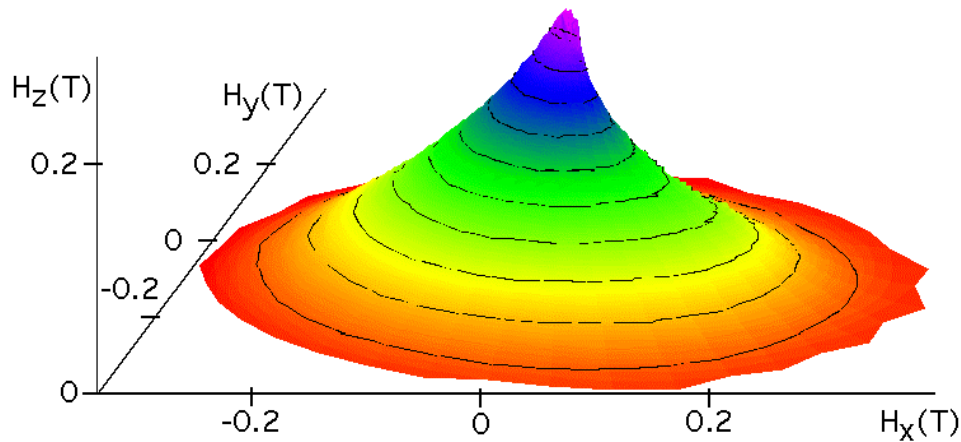
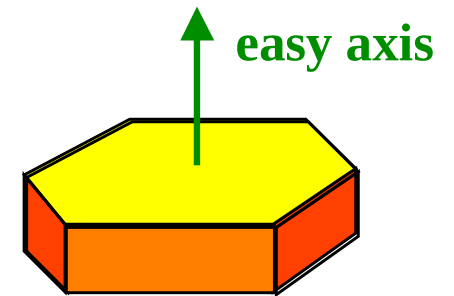
molecules



spin



Macroscopic Quantum Tunneling of Magnetization of Single Nanoparticles

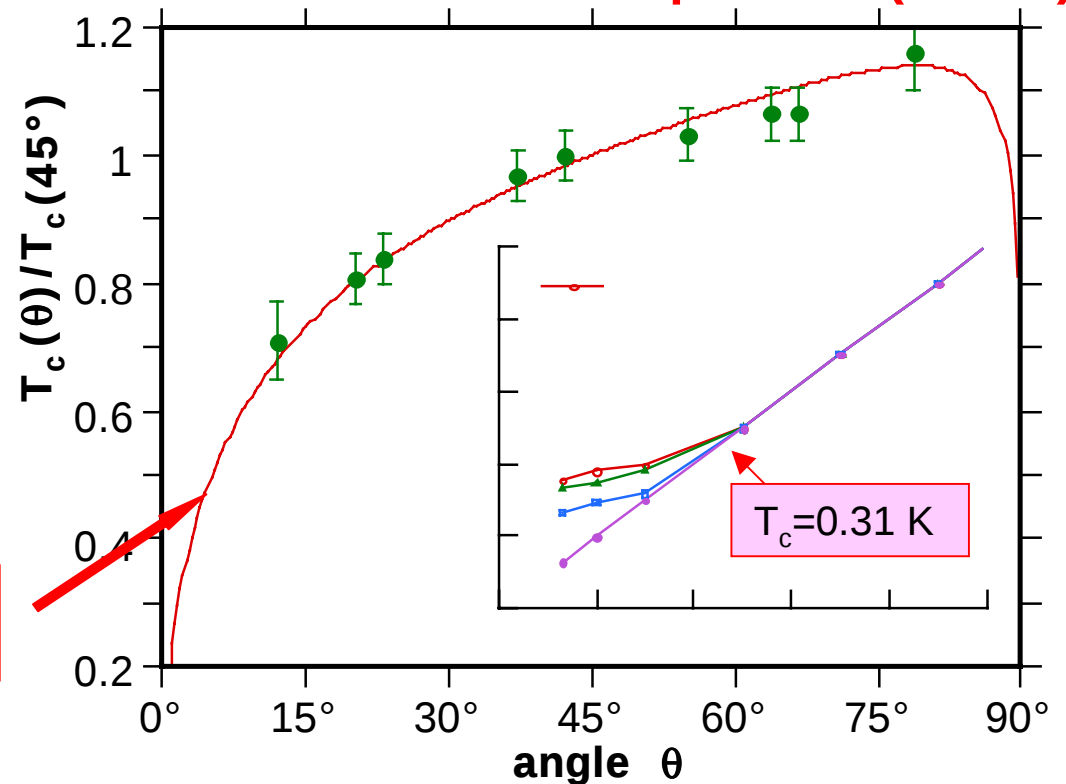


Stoner-Wohlfarth astroid

**Barium ferrite
Nanoparticle (10 nm)**

$$T_c(\theta) \propto \mu_0 H_a \epsilon^{1/4} |\cot \theta|^{1/6} (1 + |\cot \theta|^{2/3})^{-1}$$

Miguel and Chudnovsky, PRB (1995)



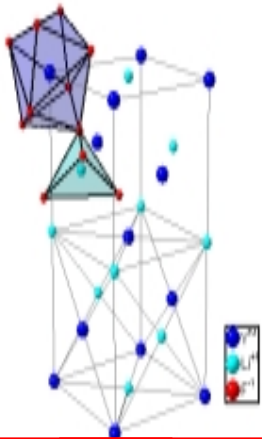
Wernsdorfer et al. et al, PRL, 79, 4014, (1997)

Conclusion and Perspectives

Quantum Tunneling at the Mesoscopic Scale
(Environmental Effects on Quantum Mechanics)

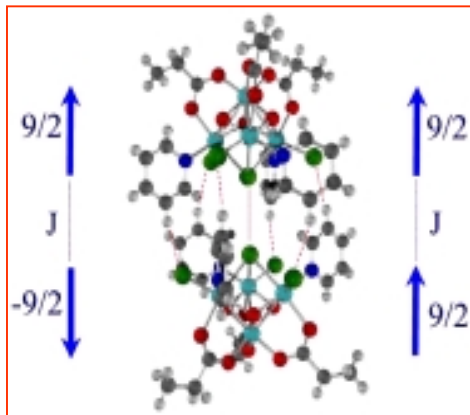
Evidence for Quantum Coherence
(τ_ϕ , Rabi oscillations, ...)

Manipulations of Quantum Spins, Spins Qbits
(Quantum Informations and Computers)



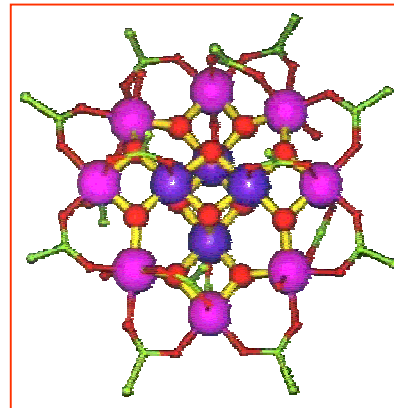
Ho³⁺

Tunneling of single Ho³⁺ ions
Entangled I-J states



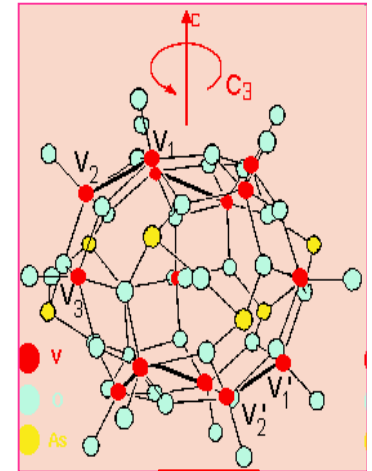
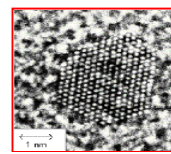
Mn₄

Ho³⁺ , Mn₄ pairs:
Cross-spin transitions, Co-tunneling



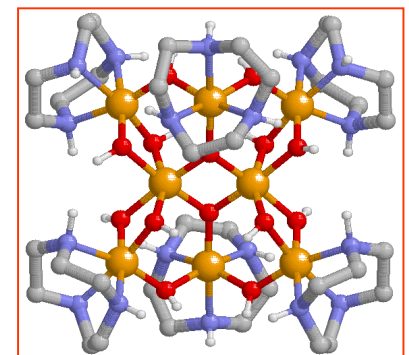
Mn₁₂-ac

Quantum → Classical crossover
Quantum Dynamics, Spin Bath



V₁₅

Dissipation control of LZS
Molecule spins 1/2 : Gapped



Fe₈

Quantum Dynamics
Berry Phases

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